



第六届全国冷原子物理与
量子信息青年学者学术讨论会

Experimental realization of spin-orbit coupling in degenerate Fermi gas

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山西大學光電研究所





Outline

Motivation: Quantum simulation with ultracold atoms

Experimental realization of spin-orbit coupling in degenerate Fermi gas

- **Raman Rabi oscillation**
- **Momentum distribution asymmetry**
- **Topological change of Fermi surface**

Momentum-resolved RF spectroscopy of spin-orbit coupling degenerate Fermi gas

Spin-orbit coupling Feshbach molecules



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Quantum simulation with ultracold atoms

We can control the Hamiltonian of cold atoms in a number of ways

$$\hat{H} = \frac{\mathbf{p}^2}{2m} + V(\mathbf{x}) + U_{\text{interaction}}$$

Kinetic: Synthetic vector potential

Potential: Optical lattice

Interaction: Feshbach resonance

For a particle with charge q , moving in an electromagnetic field, the Hamiltonian can be expressed as:

$$\hat{H} = \frac{(\mathbf{p} - q\mathbf{A})^2}{2m} + V(\mathbf{x})$$

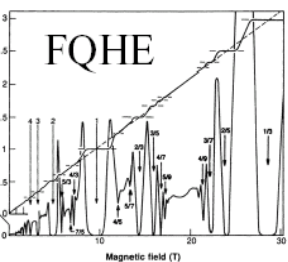
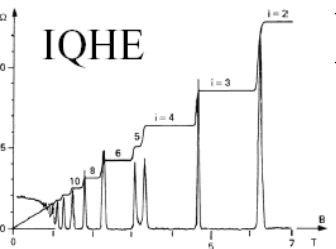
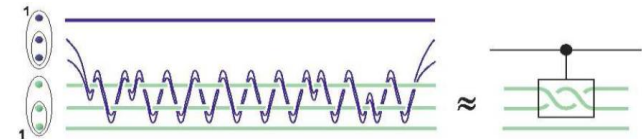
Vector potential: \mathbf{A}

Magnetic field: $\mathbf{B} = \nabla \times \mathbf{A}$

Once we could construct such a Hamiltonian for the neutral atoms, we can simulate the charged particle with neutral atoms!!

Ultracold atoms simulate charged particle

- To simulate Lorenz Force: $F = q\vec{v} \times \vec{B}$
- To understand Quantum Hall Effect?
- To form the topological insulator
- Topological Quantum Computing



Klaus von Klitzing

Nobel Prize 1985
Quantum Hall Effect



Robert B. Laughlin



Horst L. Störmer



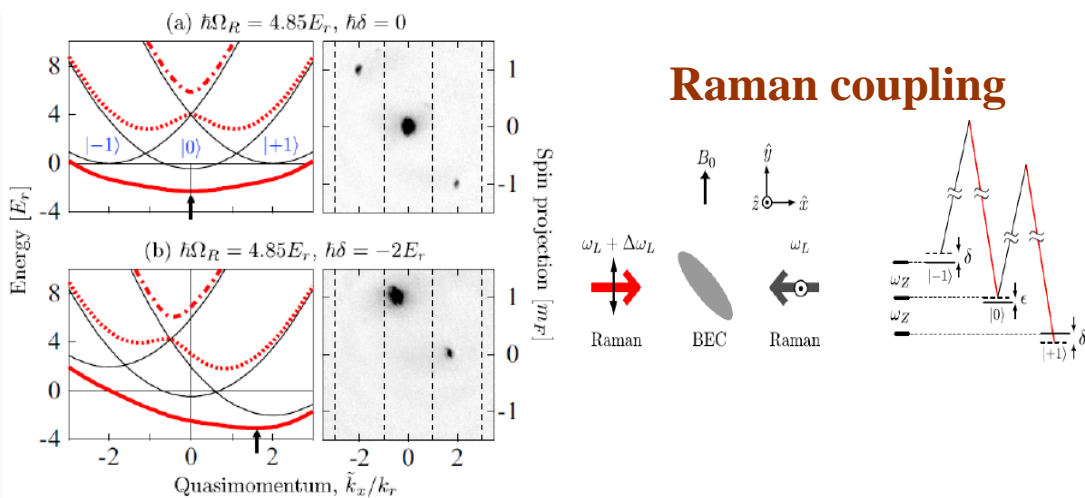
Daniel C. Tsui

Nobel Prize 1998
Fractional Quantum Hall Effect

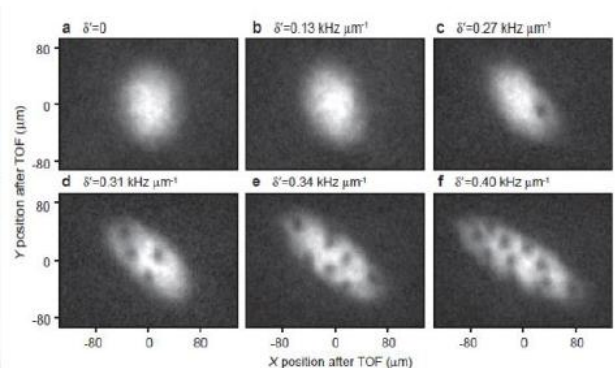
Experimental Progress

Synthetic magnetic fields for BEC

Effective Vector Potential generation:



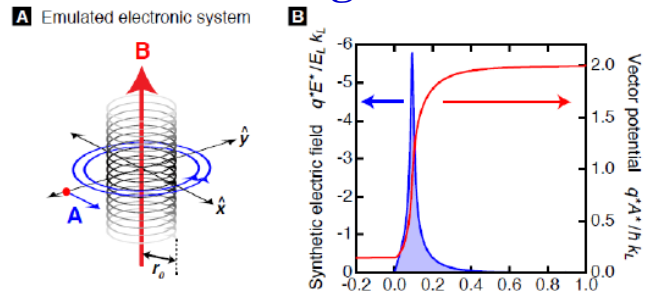
Y.-J. Lin, et.al., PRL 102, 130401 (2009)



Vortices are formed in condensates

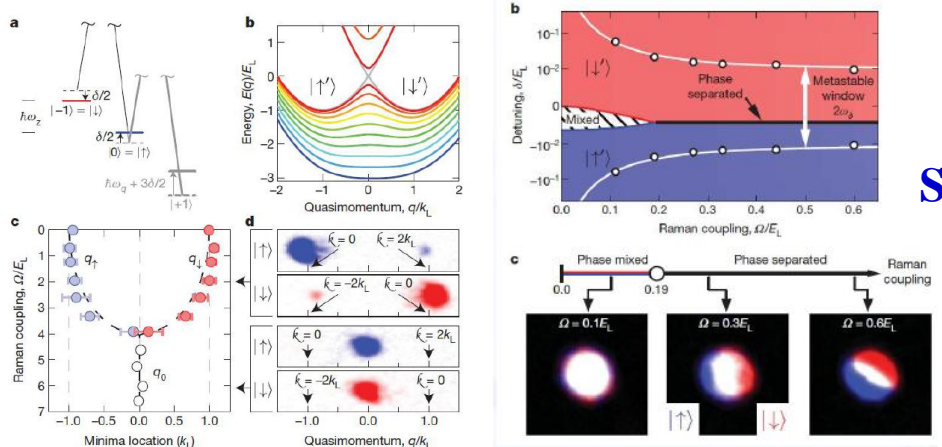
Y.-J. Lin, et.al., Nature 462, 628 (2009)

Electric field generation



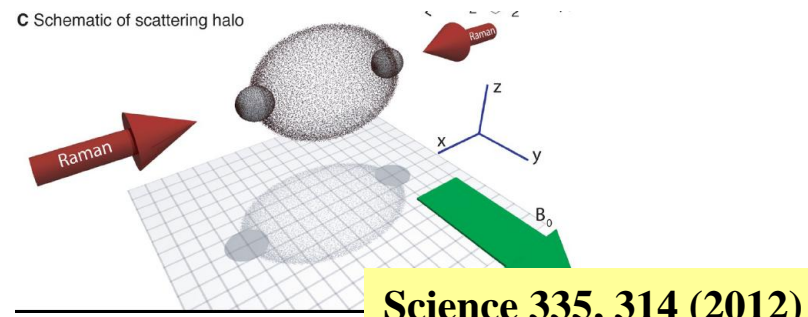
Y.-J. Lin, et.al., Nature Physics 7, 531 (2011)

Spin-Orbit coupling



Y. -J. Lin, et.al., Nature 417, 83 (2011)

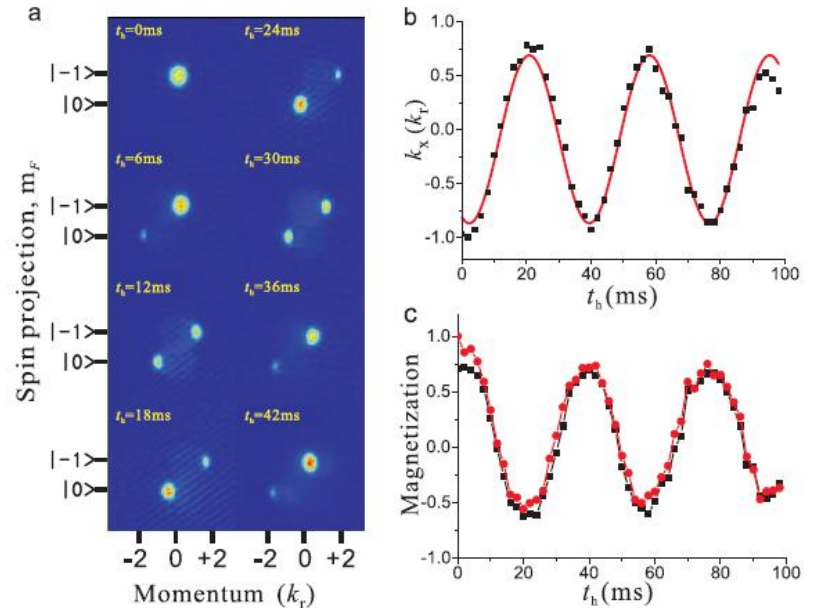
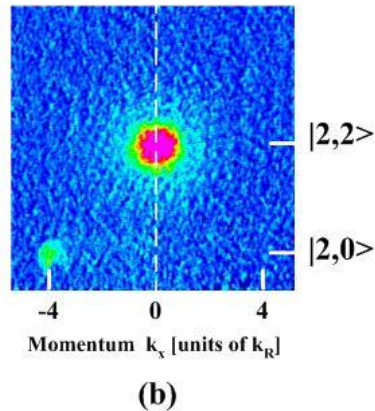
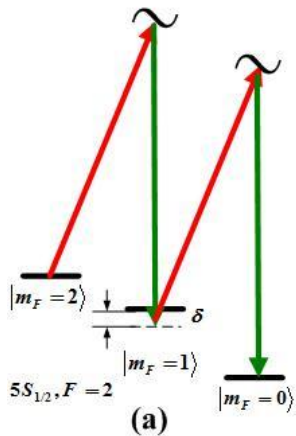
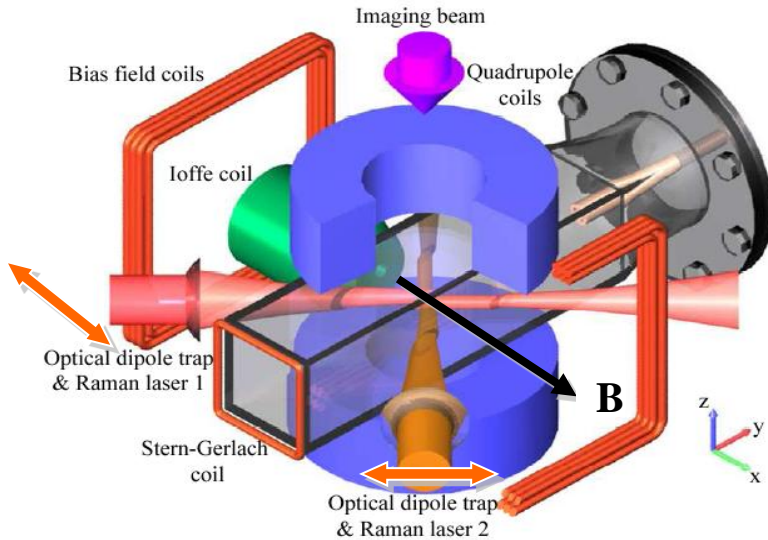
Synthetic Partial Waves in Atomic Collisions



Science 335, 314 (2012)

BEC in light-induced vector gauge potential using the 1064 nm optical dipole trap lasers

Collective Dipole Oscillation of spin-orbit coupling BEC

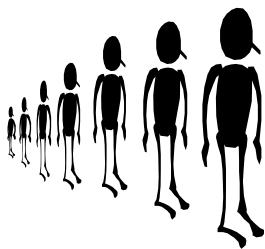


University of Science and Technology of China,
Shuai Chen and Jian-Wei Pan's group

Z. Fu, P. Wang, S. Chai, L. Huang,
J. Zhang, *Phys. Rev. A* 84, 043609 (2011)

S. Chen, J. Y. Zhang, S. C. Ji, Z. Chen, L.
Zhang, Z. D. Du, Y. J. Deng, H. Zhai, and J.
W. Pan, arXiv:1201.6018.

超冷原子气体中的合成规范场



Bosons

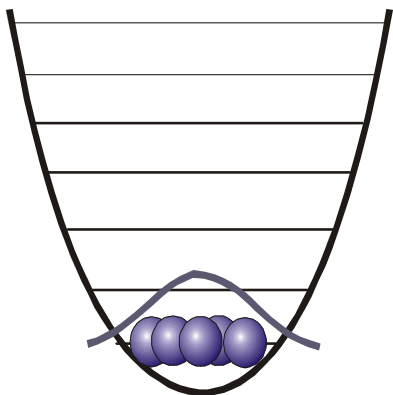
- Integer spin



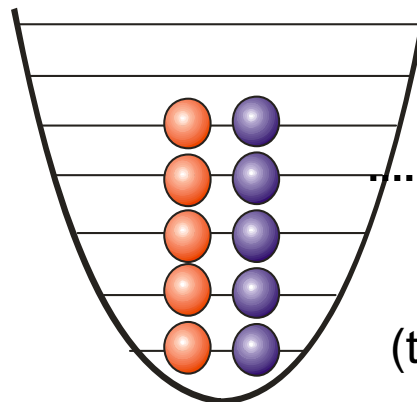
Fermions

- Half-integer spin

Atoms in a harmonic potential.



$T = 0$



$E_F = k_b T_F$

(two spin states)

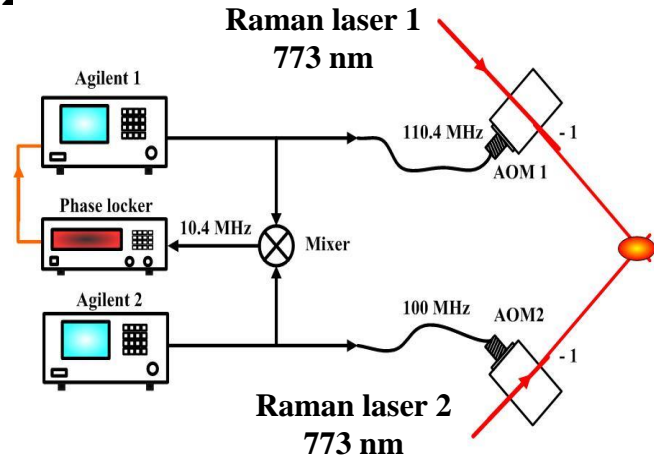
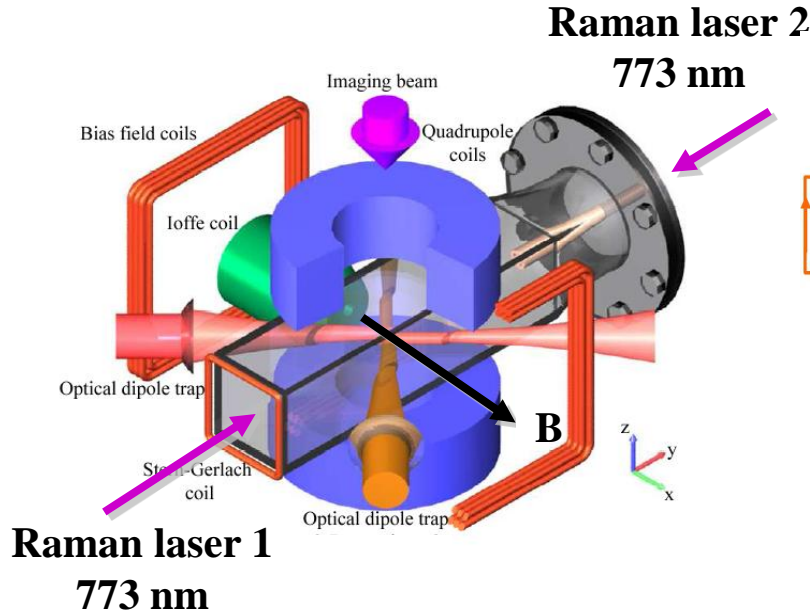
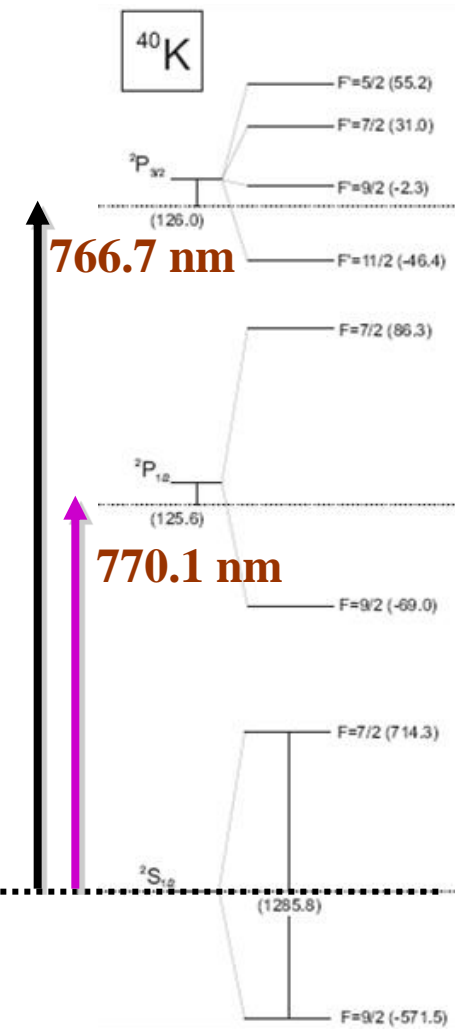
Bose-Einstein condensation

NIST, 山西大学, 科大

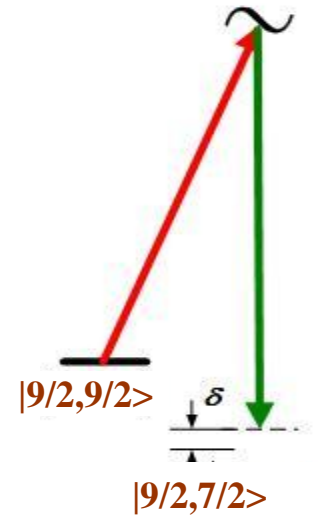
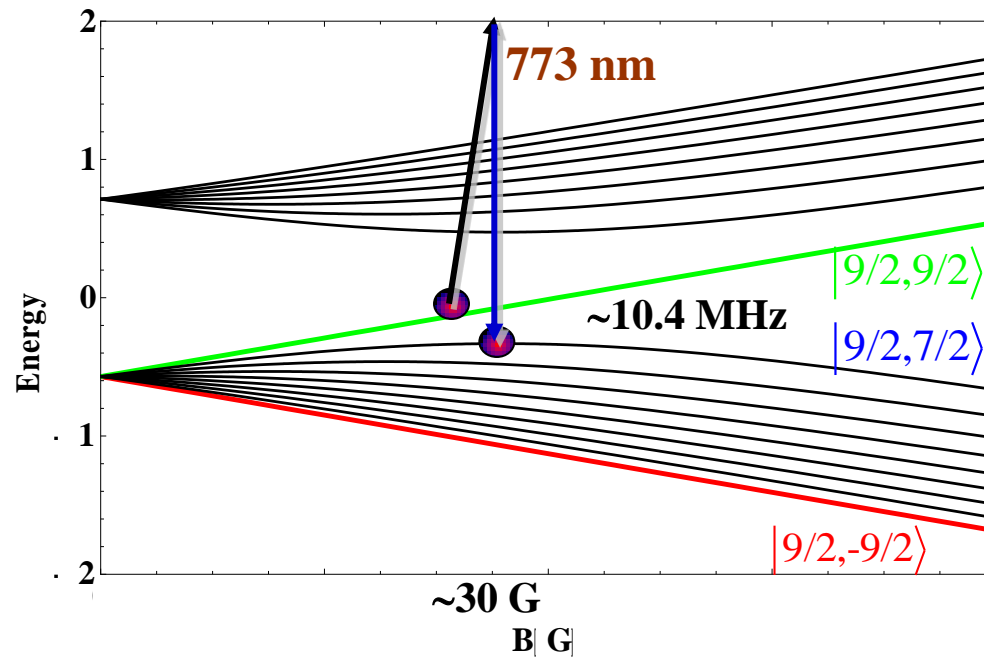
Degenerate Fermi gas

山西大学, MIT

Experimental realization of spin-orbit coupling in degenerate Fermi gas



Raman laser 1
773 nm



$|9/2, 7/2\rangle \rightarrow |9/2, 5/2\rangle$
 $\delta \sim 170\text{ KHz}$

Theoretical model of two-level system

$$\mathbf{H} = \frac{\hbar^2 \mathbf{p}^2}{2m} \hat{\mathbf{I}} + \frac{\delta}{2} \hat{\sigma}_z + \frac{\Omega}{2} \hat{\sigma}_x \cos(2\mathbf{k}_R \hat{\mathbf{x}}) - \frac{\Omega}{2} \hat{\sigma}_y \sin(2\mathbf{k}_R \hat{\mathbf{x}}) \quad (\text{NIST group's SO Coupling})$$

$$U = \begin{pmatrix} e^{-ik_R x} & 0 \\ 0 & e^{ik_R x} \end{pmatrix}$$



Translate unitary transformation

Base: $\{ |\uparrow, k_x = p + k_R\rangle, |\downarrow, k_x = p - k_R\rangle \}$

$$H_R(k_x) = \hbar \begin{pmatrix} \frac{\hbar}{2m}(p + k_R)^2 - \delta/2 & \Omega/2 \\ \Omega/2 & \frac{\hbar}{2m}(p - k_R)^2 + \delta/2 \end{pmatrix}$$



two energy eigenvalues:

$$E_{\pm}(p) = \hbar \left[\frac{\hbar(p^2 + k_R^2)}{2m} \pm \sqrt{(4\hbar p k_R / 2m - \delta)^2 + \Omega^2 / 2} \right]$$

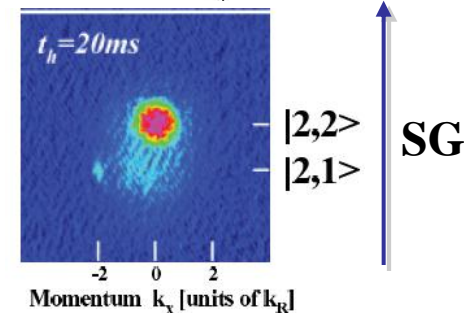
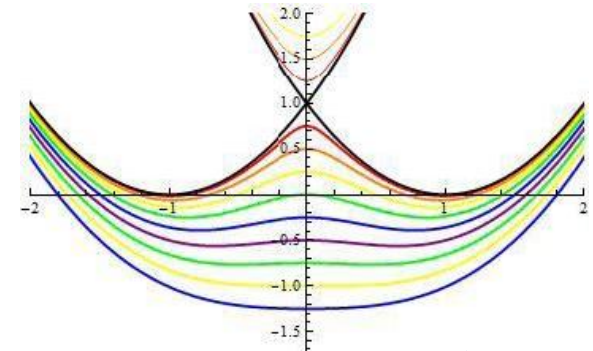
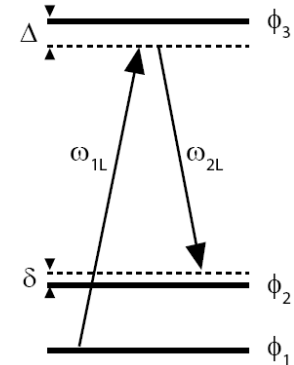
two dressed eigenstates:

$$|\uparrow', p\rangle = c_1 |\uparrow, k_x = p + k_R\rangle + c_2 |\downarrow, k_x = p - k_R\rangle$$

$$|\downarrow', p\rangle = c_3 |\uparrow, k_x = p + k_R\rangle + c_4 |\downarrow, k_x = p - k_R\rangle$$

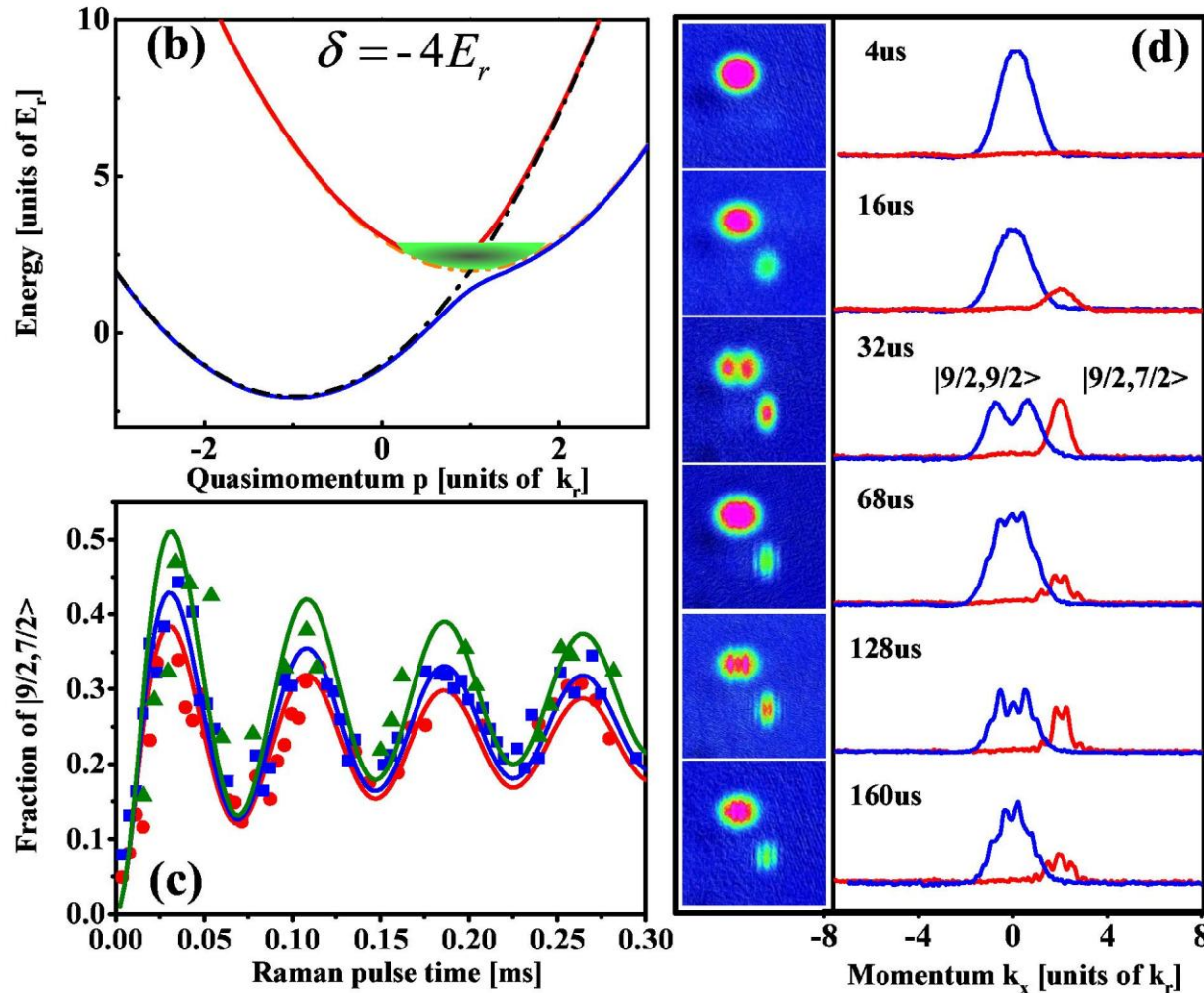
$$\hat{H}_2 = \frac{\hbar^2 \hat{\mathbf{k}}^2}{2m} \check{\mathbf{I}} + \frac{\Omega}{2} \check{\sigma}_z + \frac{\delta}{2} \check{\sigma}_y + 2 \frac{\hbar^2 k_L \hat{k}_x}{2m} \check{\sigma}_y + E_L \check{\mathbf{I}}$$

Spin-orbit coupled form

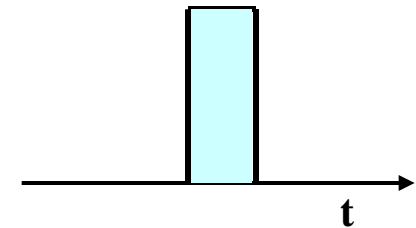


Raman laser

SO Coupled Fermi Gases: Raman Rabi Oscillation



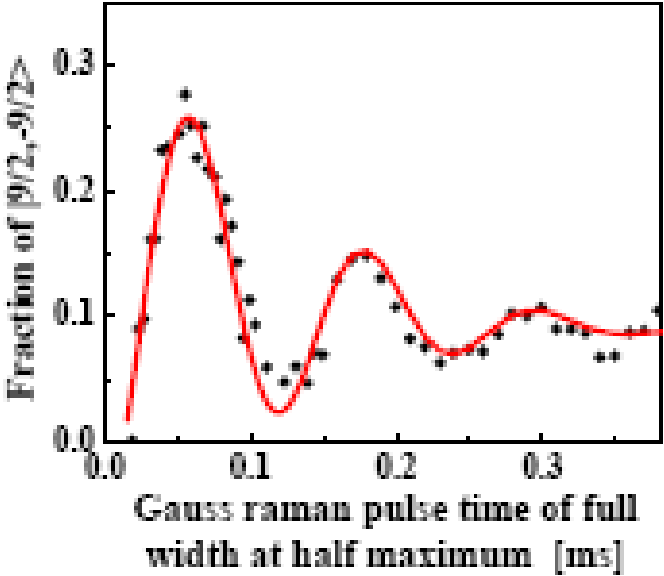
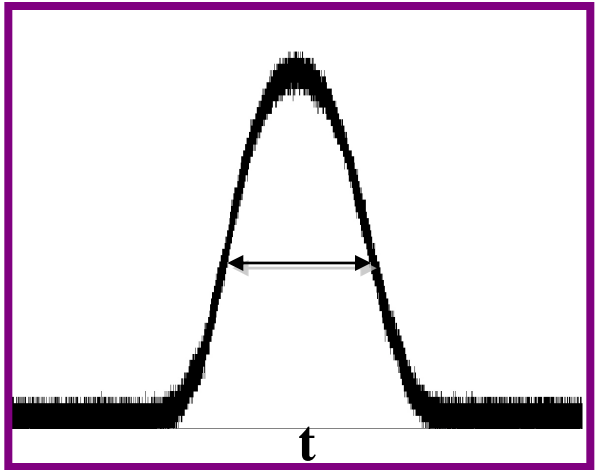
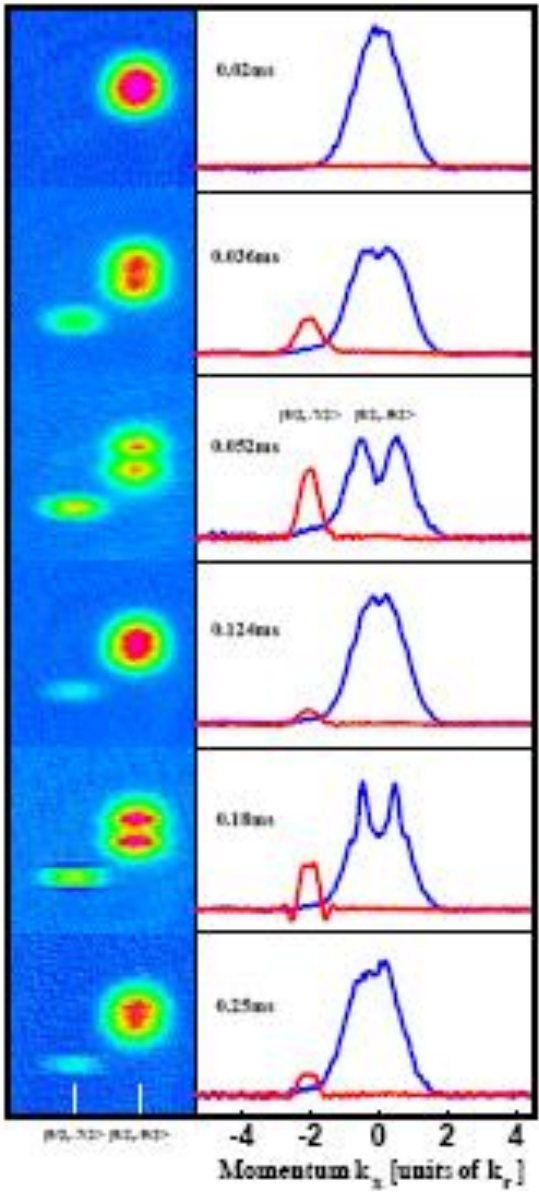
First prepare fermion in $9/2$, and then turn on Raman coupling with square envelop pulse



$$n_{\uparrow}(\mathbf{p} + 2\hbar k_r \hat{e}_z, \mathbf{r}, t) = n_{\uparrow}(\mathbf{p}, \mathbf{r}, 0) \frac{1}{1 + \left(\frac{2p_x k_r}{\Omega m}\right)^2} \sin^2 \sqrt{(p_x k_r / m)^2 + \Omega^2 / 4t}$$

$$\Omega = 1.48(5) E_r.$$

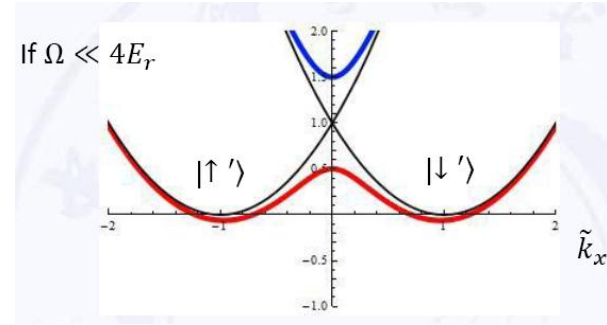
Raman coupling with Gaussian envelop pulse



SO Coupled Fermi Gases: Equilibrium Momentum distribution

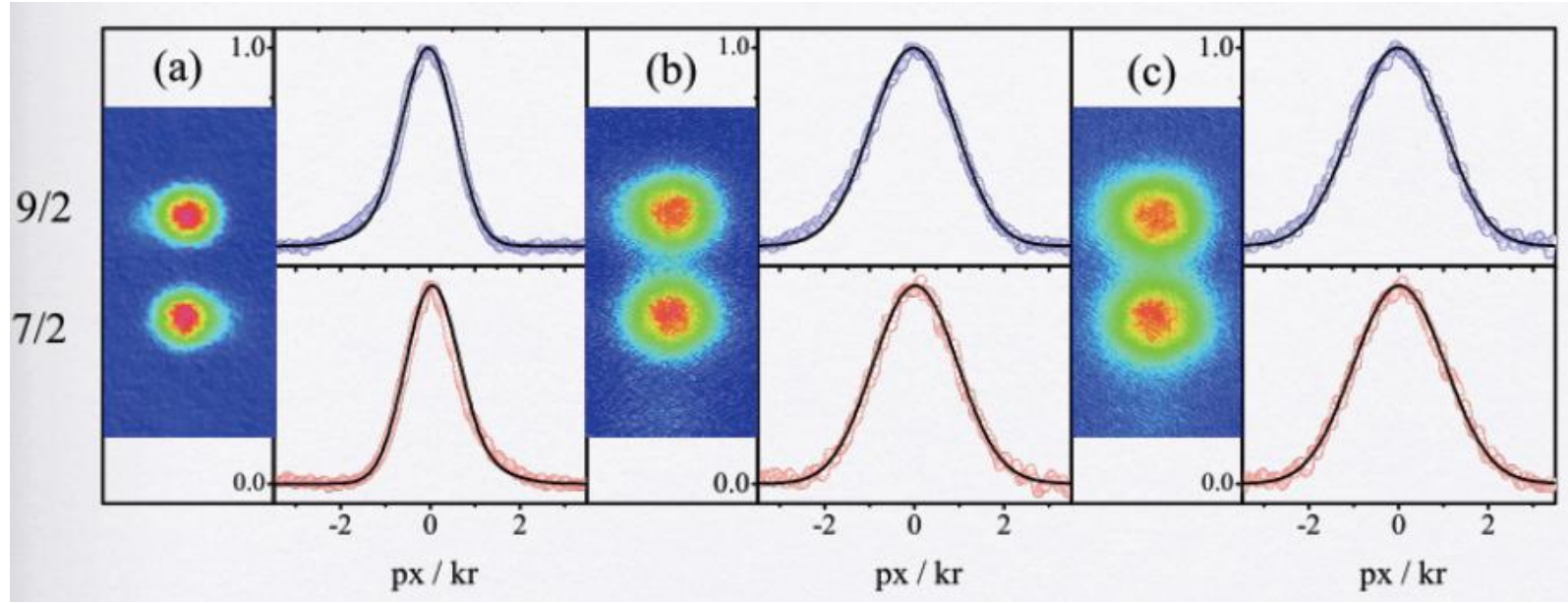
$$H_R(k_x) = \hbar \begin{pmatrix} \frac{\hbar}{2m}(p + k_R)^2 - \delta/2 & \Omega/2 \\ \Omega/2 & \frac{\hbar}{2m}(p - k_R)^2 + \delta/2 \end{pmatrix}$$

$$\delta = 0.$$



Break spatial reflectional symmetry: $n(\mathbf{p}) \neq n(-\mathbf{p})$

Preserve reversal symmetry: $n_{\uparrow}(\mathbf{k}) = n_{\downarrow}(-\mathbf{k})$

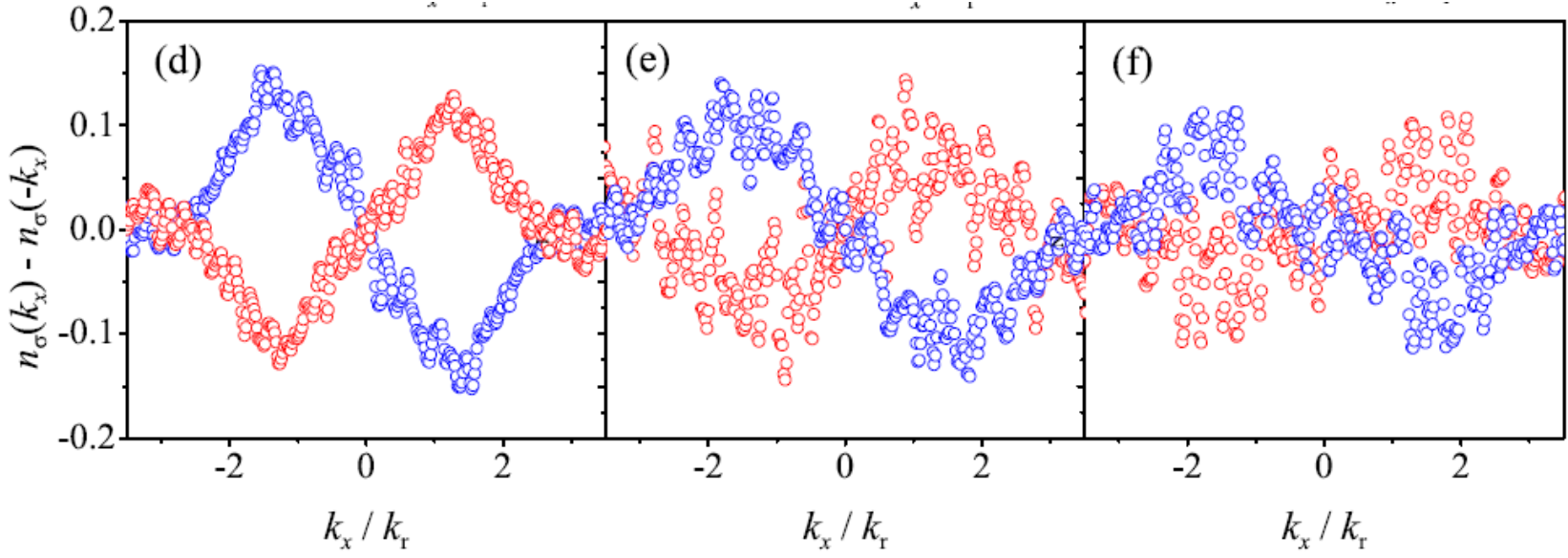
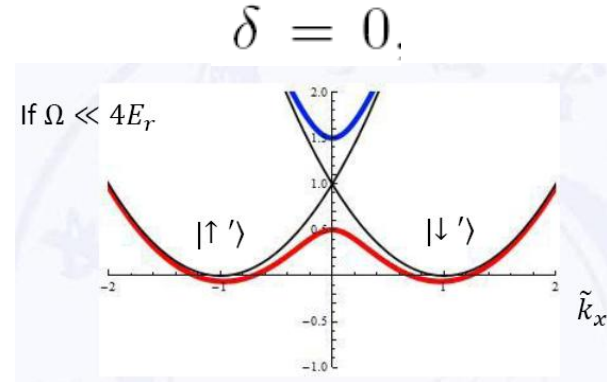


Time of flight measurement with Stern-Gerlach effect

SO Coupled Fermi Gases: Equilibrium Momentum distribution

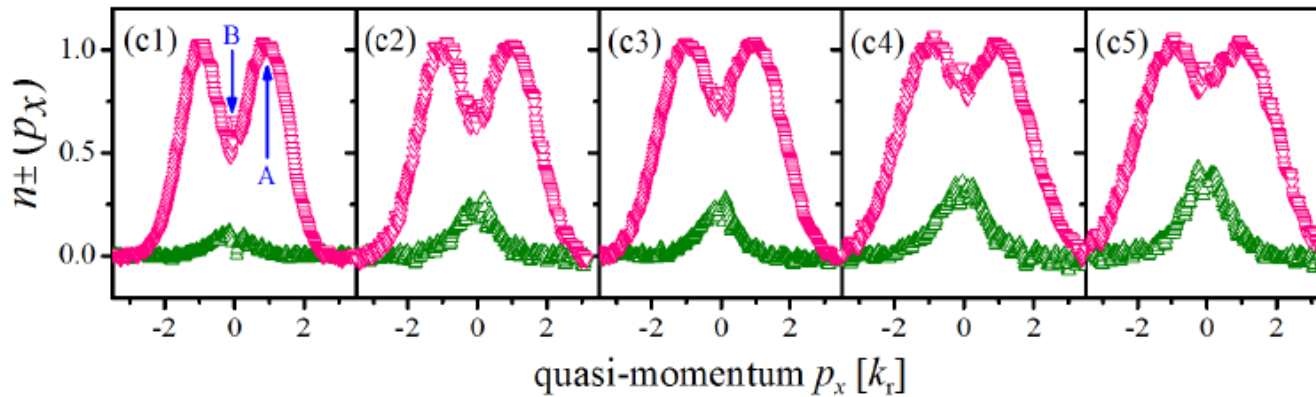
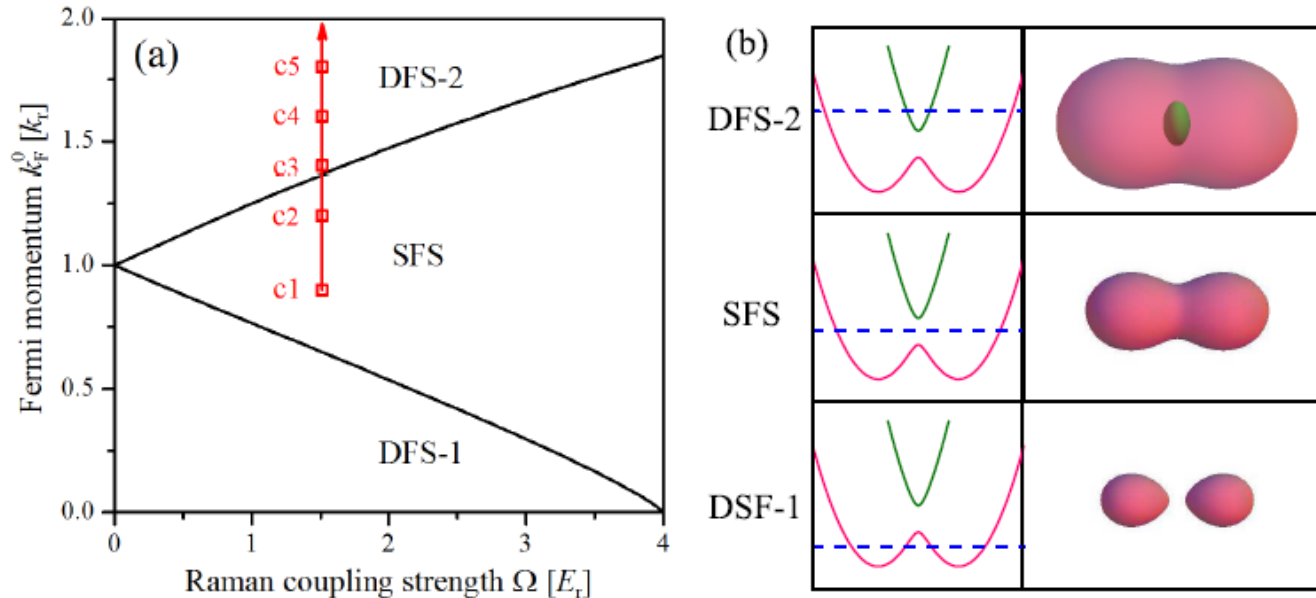
Break spatial reflectional symmetry: $n(\mathbf{p}) \neq n(-\mathbf{p})$

Preserve time reversal symmetry: $n_{\uparrow}(\mathbf{k}) = n_{\downarrow}(-\mathbf{k})$



$$n_{\sigma}(p_x) - n_{\sigma}(-p_x)$$

SO Coupled Fermi Gases: Momentum distribution in helical bases



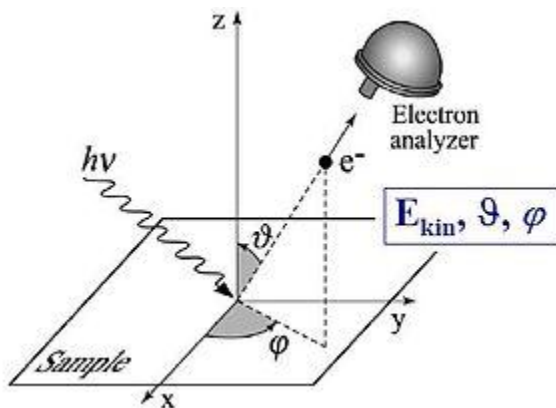
$$n_+(\mathbf{p}) = \frac{u_{\mathbf{p}}^2 n_{\uparrow}(\mathbf{p} - \hbar k_r \hat{\mathbf{e}}_x) - v_{\mathbf{p}}^2 n_{\downarrow}(\mathbf{p} + \hbar k_r \hat{\mathbf{e}}_x)}{u_{\mathbf{p}}^2 - v_{\mathbf{p}}^2}$$

$$n_-(\mathbf{p}) = \frac{v_{\mathbf{p}}^2 n_{\uparrow}(\mathbf{p} - \hbar k_r \hat{\mathbf{e}}_x) - u_{\mathbf{p}}^2 n_{\downarrow}(\mathbf{p} + \hbar k_r \hat{\mathbf{e}}_x)}{v_{\mathbf{p}}^2 - u_{\mathbf{p}}^2}$$

(A) Double peak structure in lower branch
(B) Small population in higher branch

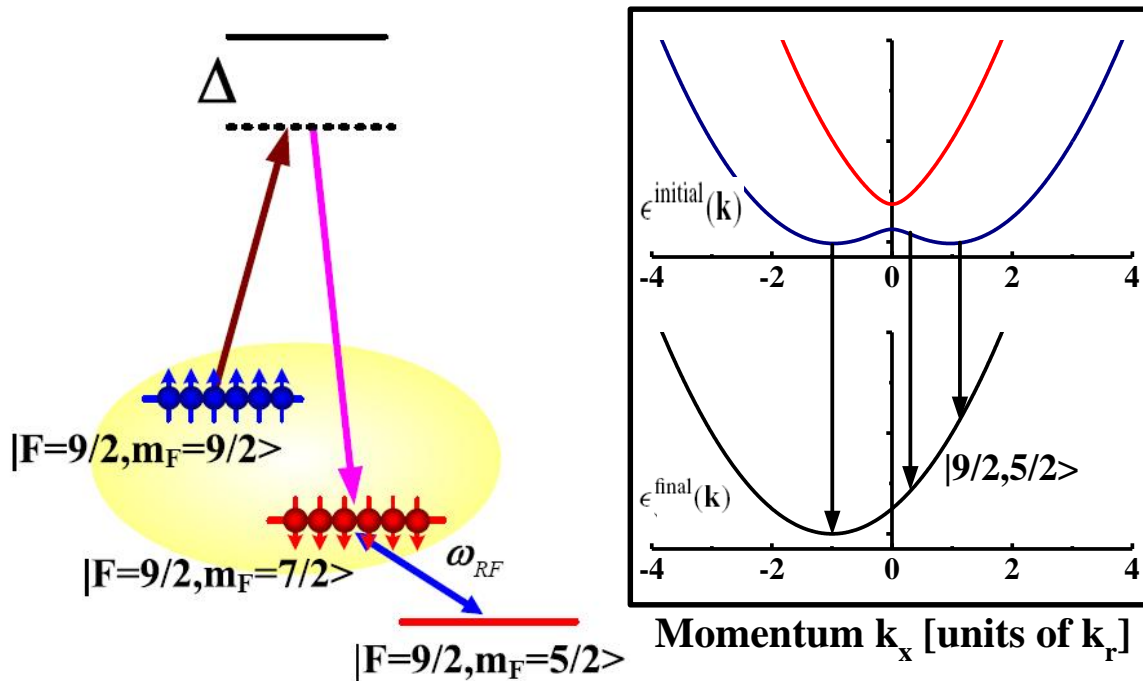
(A) Double peak structure gradually disappears
(B) Significant population in higher branch

Momentum-resolved RF spectroscopy of non-interacting SO coupling Fermi gas



角分辨光电子谱 (Angle resolved photoemission spectroscopy ARPES)

coupling Fermi gas



$$\hbar\omega_{RF} = E_Z + \epsilon^{\text{initial}}(\mathbf{k}) - \epsilon^{\text{final}}(\mathbf{k})$$

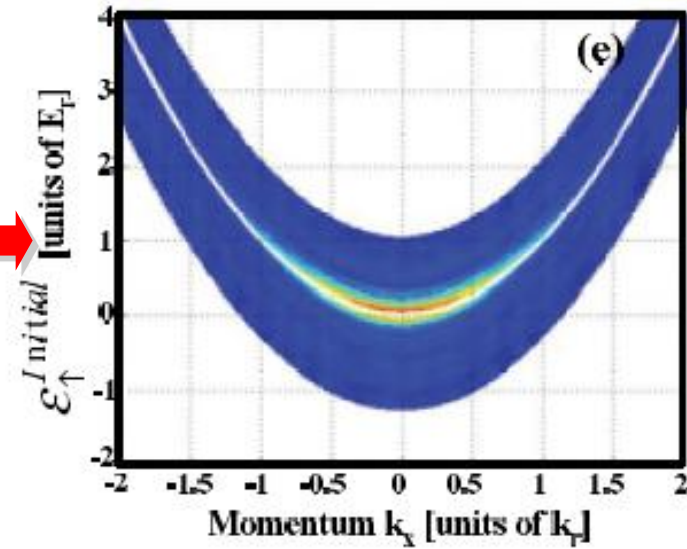
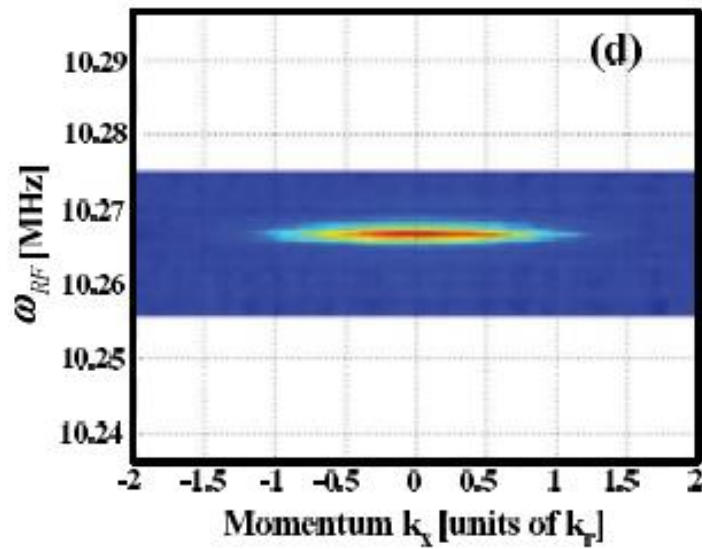
➔
$$\epsilon^{\text{initial}}(\mathbf{k}) = \hbar\omega_{RF} - E_Z + \epsilon^{\text{final}}(\mathbf{k})$$

E_Z : the energy split of the two Zeeman states

$\epsilon^{\text{initial}}(\mathbf{k})$: energy-momentum dispersion of the initial state

$\epsilon^{\text{final}}(\mathbf{k})$: energy-momentum dispersion of the final state (empty state)

Momentum-resolved RF spectroscopy of non-interacting Fermi gas without SO coupling



When: $\epsilon^{\text{initial}}(\mathbf{k}) = \epsilon^{\text{final}}(\mathbf{k}) = \hbar^2 \mathbf{k}^2 / 2m$

$$\hbar \omega_{\text{RF}} = E_Z$$

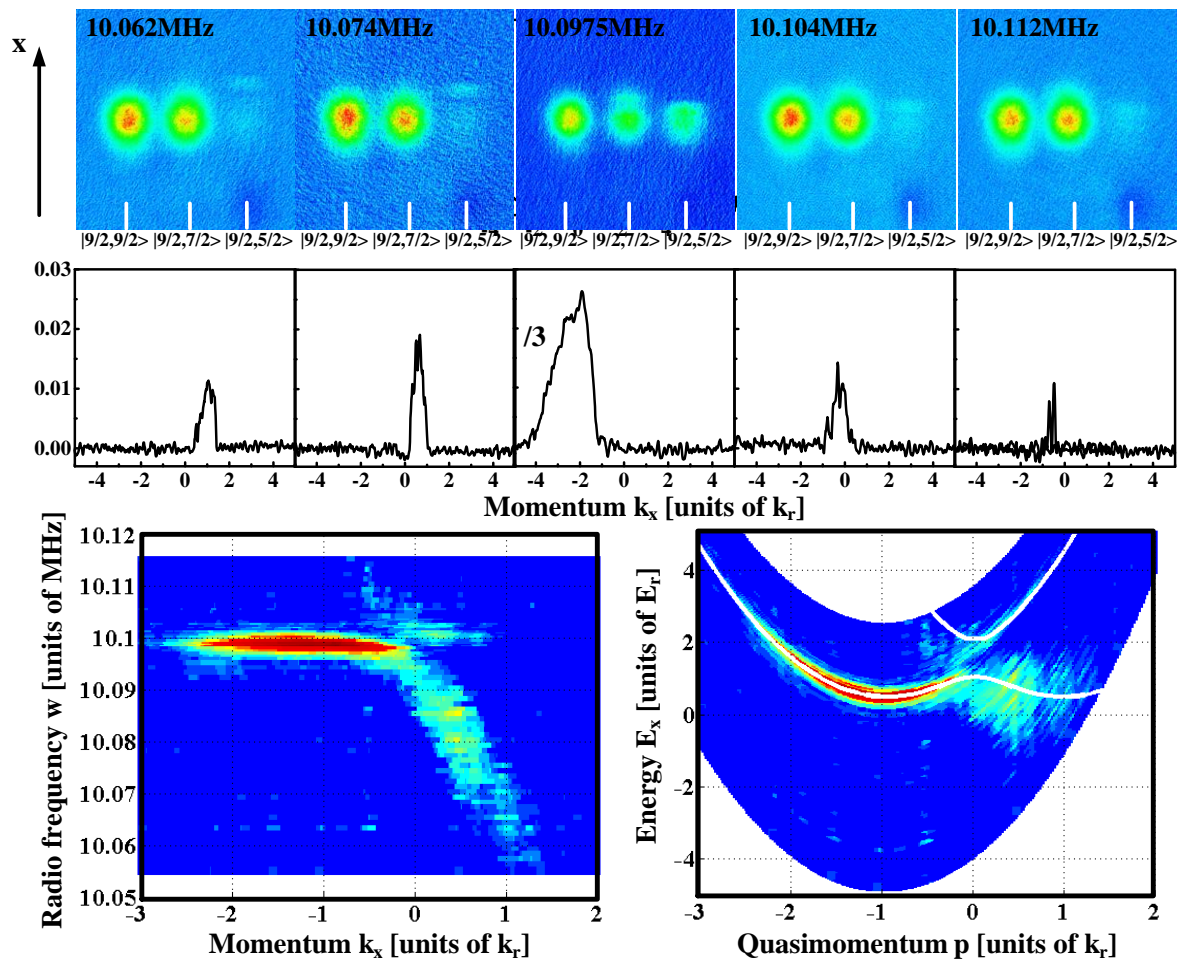
When we know:

$$\epsilon^{\text{final}}(\mathbf{k}) = \hbar^2 \mathbf{k}^2 / 2m$$

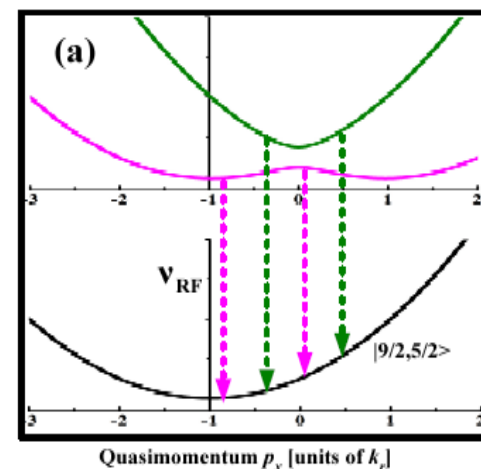
Then:

$$\epsilon^{\text{initial}}(\mathbf{k}) = \hbar \omega_{\text{RF}} - E_Z + \epsilon^{\text{final}}(\mathbf{k})$$

Momentum-resolved RF spectroscopy of non-interacting SO coupling Fermi gas



$$\delta=0 \quad \Omega=1.5E_r$$



When we know:

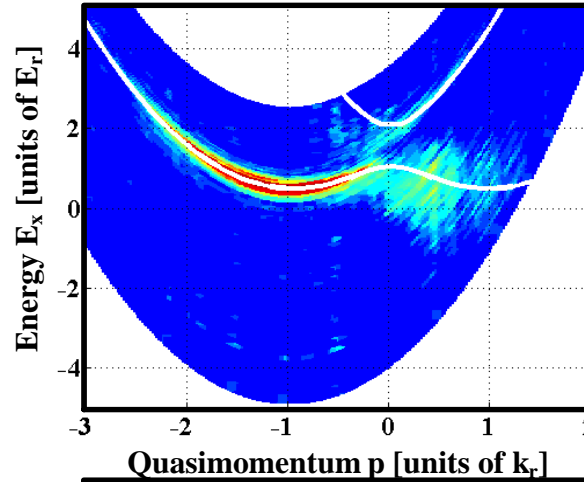
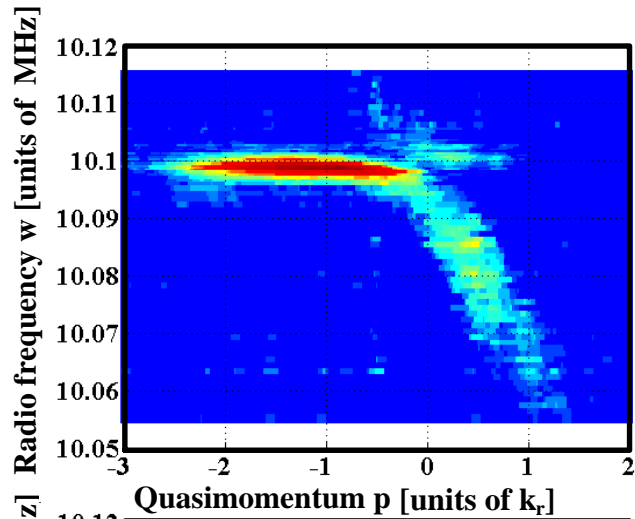
$$\epsilon^{\text{final}}(\mathbf{k}) = \hbar^2 \mathbf{k}^2 / 2m$$

Then:

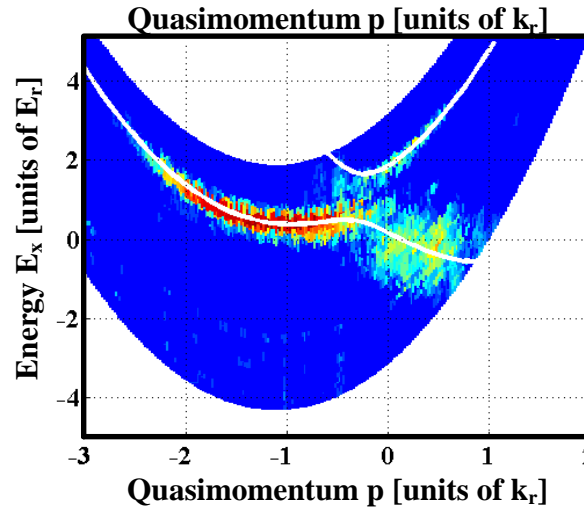
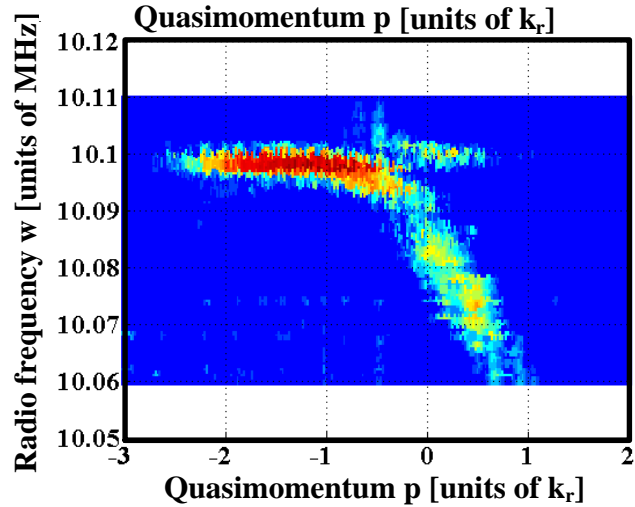
$$\epsilon^{\text{initial}}(\mathbf{k}) = \hbar\omega_{\text{RF}} - E_Z + \epsilon^{\text{final}}(\mathbf{k})$$

$$\hbar\omega_{\text{RF}} = E_Z + \epsilon_{\uparrow}^{\text{initial}}(\mathbf{k}) - \epsilon_{\downarrow}^{\text{final}}(\mathbf{k})$$

Momentum-resolved RF spectroscopy of non-interacting SO coupling Fermi gas

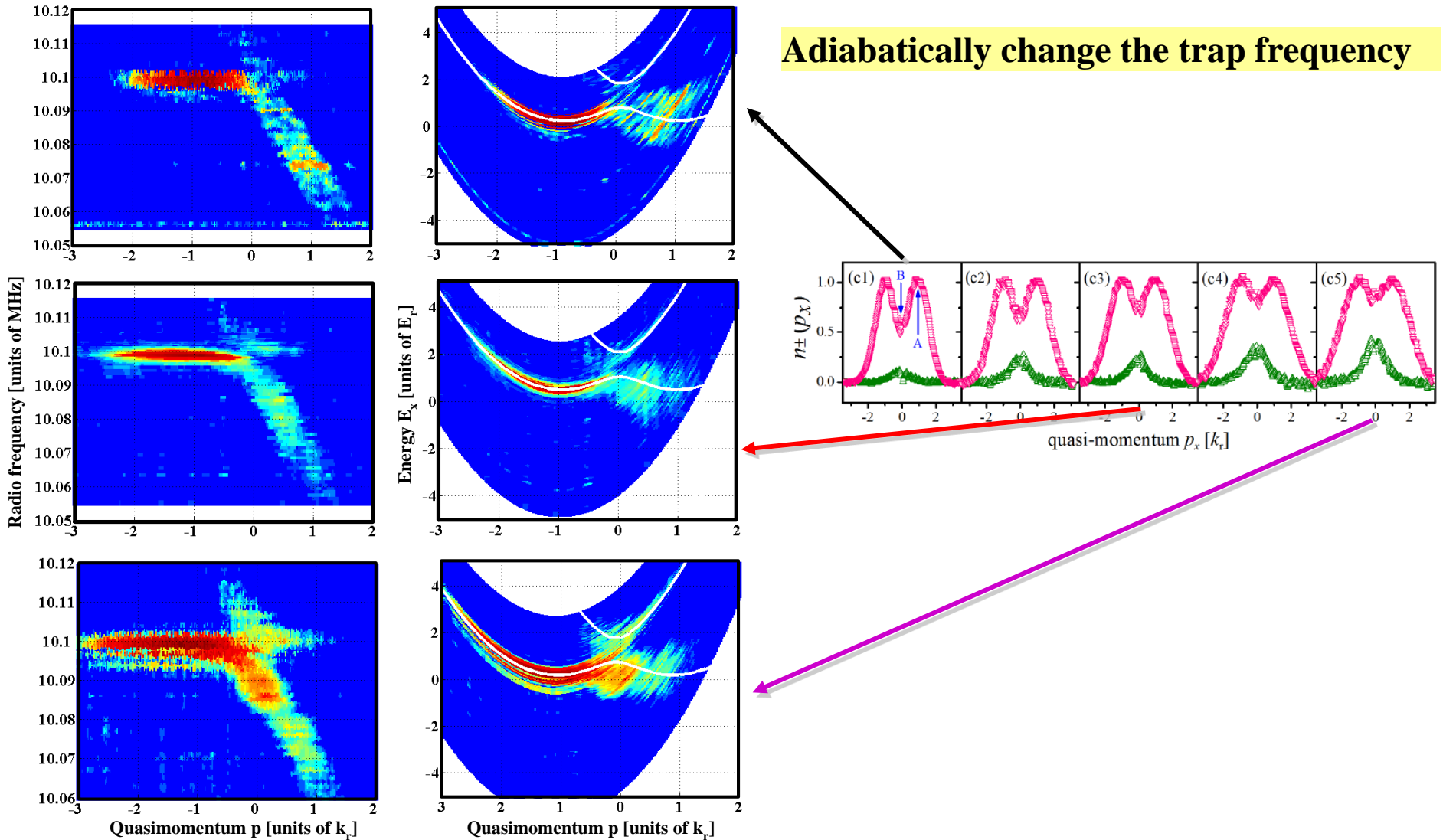


$\delta=0$



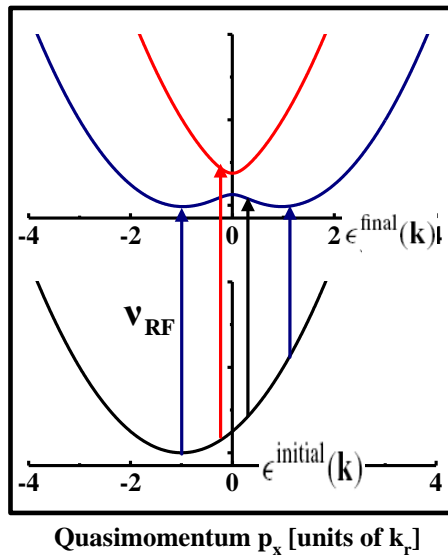
$\delta=1 E_r$

Momentum-resolved RF spectroscopy of non-interacting SO coupling Fermi gas



Spin-Injection Spectroscopy of a Spin-Orbit Coupled Fermi Gas

L. W. Cheuk, A. T. Sommer, Z. Hadzibabic, T. Yefsah, W. S. Bakr, M. W. Zwierlein, arXiv:1205.3483



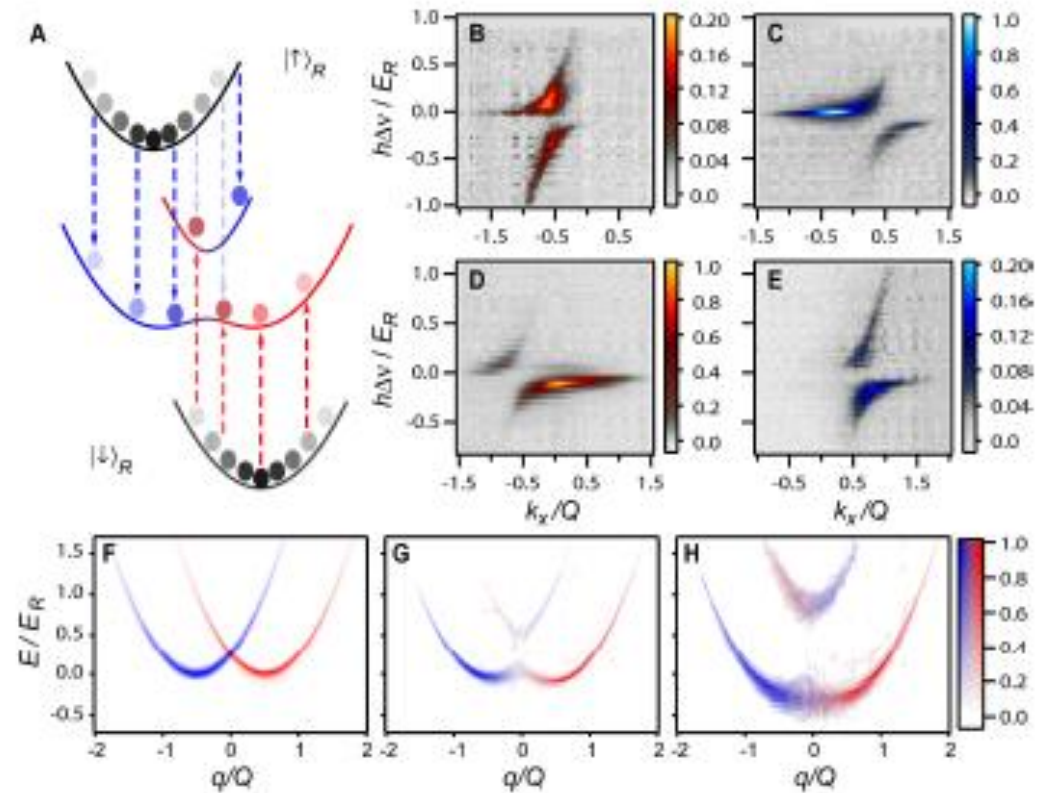
When we know:

$$\epsilon^{\text{initial}}(\mathbf{k}) = \hbar^2 \mathbf{k}^2 / 2m$$

Then:

$$\epsilon^{\text{final}}(\mathbf{k}) = \hbar\omega_{\text{RF}} - E_Z + \epsilon^{\text{initial}}(\mathbf{k})$$

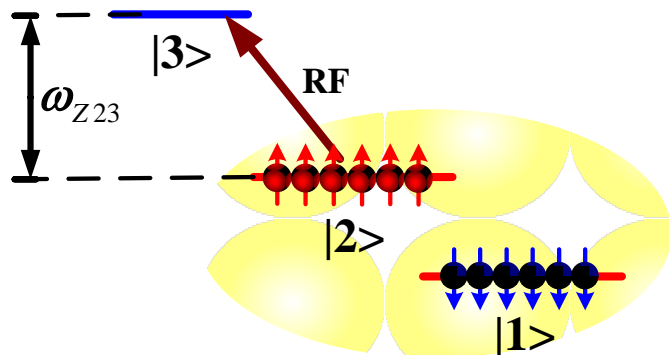
$$\epsilon^{\text{final}}(\mathbf{k}) \quad \text{:SO coupling}$$



Raman fields have previously been used to generate spin-orbit coupling and gauge fields in pioneering work on Bose-Einstein condensates [19–21], and recently spin-orbit coupling in Fermi gases [22]. Here, we directly

[22] P. Wang, *et al.* Spin-Orbit Coupled Degenerate Fermi Gases. *arXiv:1204.1887v1 [cond-mat.quant-gas]* (2012).

RF spectroscopy of strongly interacting ultracold Fermi gas

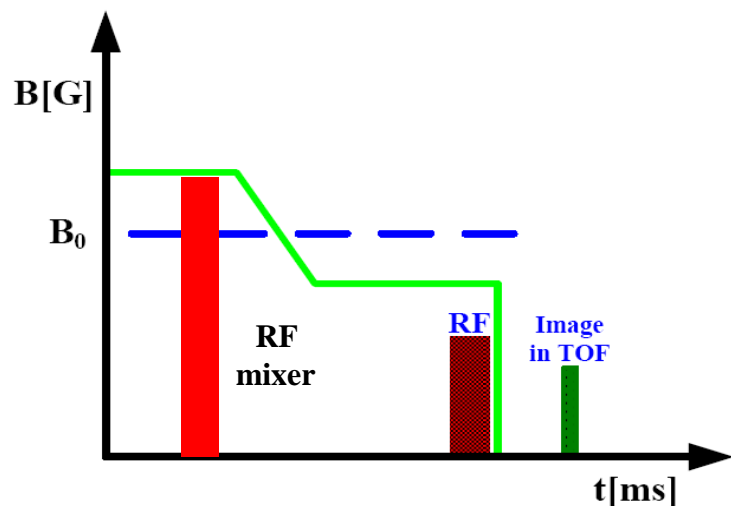


$$\text{---} \bullet \text{---} \quad |9/2, -7/2\rangle$$

$$\text{---} \bullet \text{---} \quad |9/2, -9/2\rangle$$

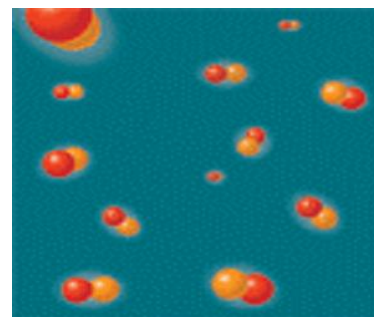
S-wave $B_0 = 202.2$ G

$$|9/2, -9/2\rangle + |9/2, -7/2\rangle$$



BEC – BCS Crossover

molecules
 $\alpha > 0$

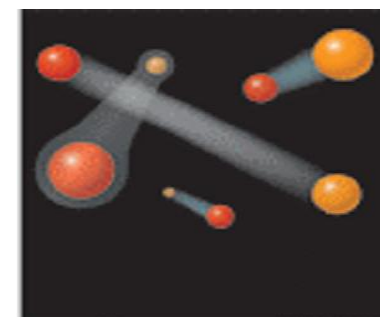


localized pairs

Crossover

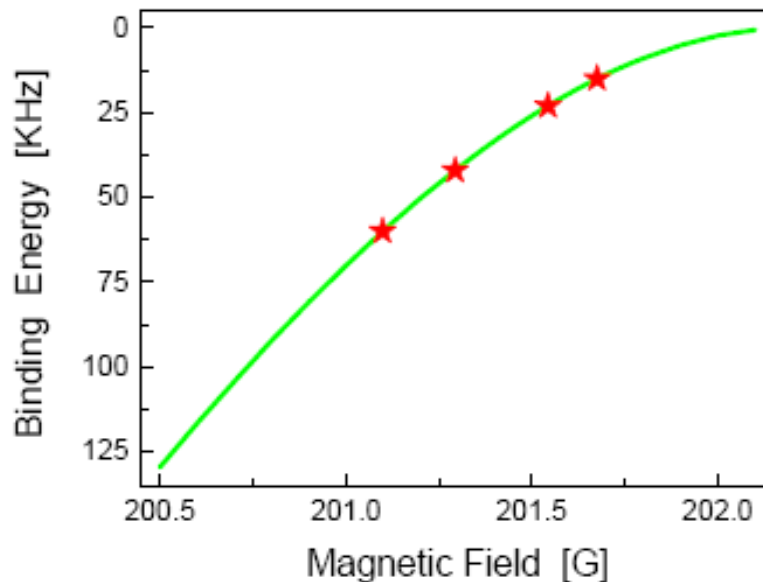
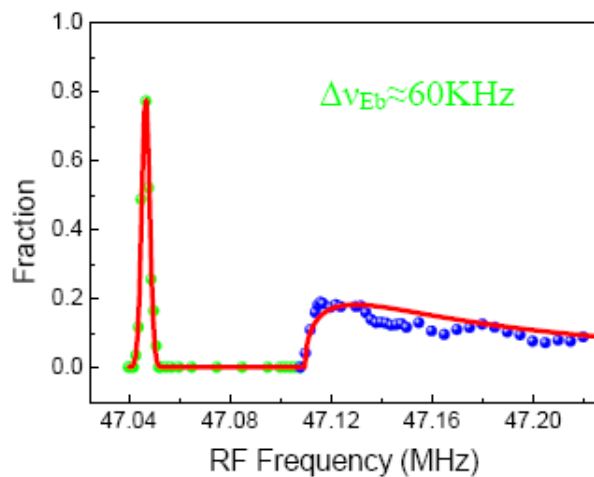
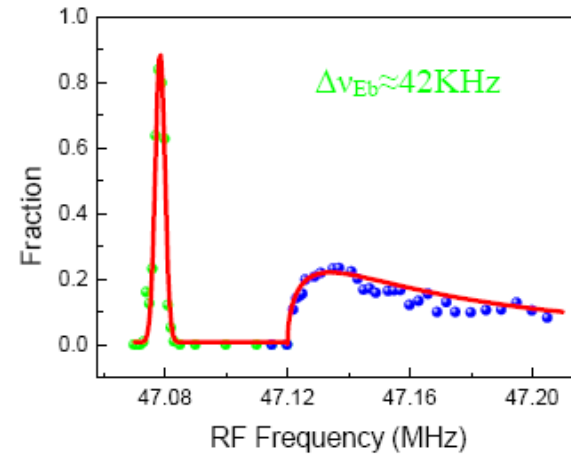
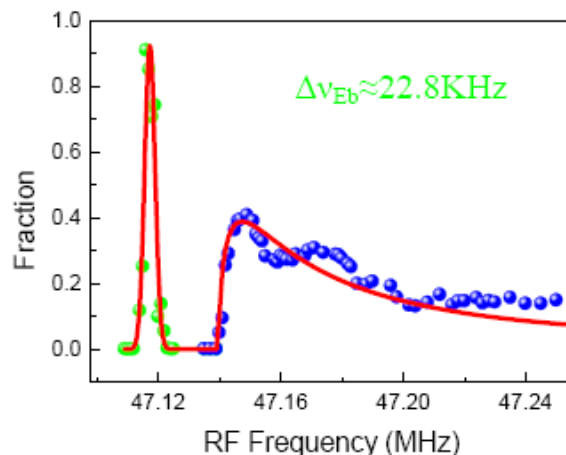
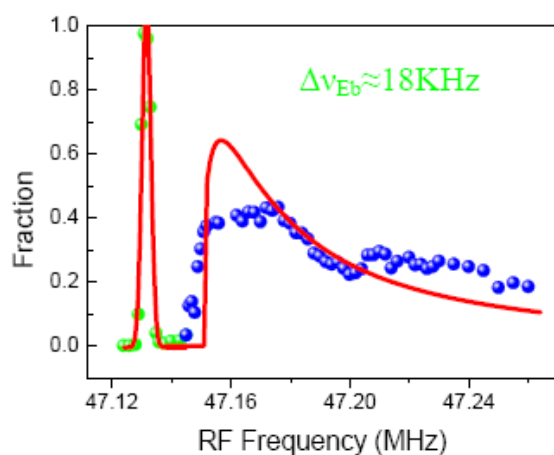
Feshbach resonance

Cooper pairs
weak coupling $\alpha < 0$



Nonlocalized pairs

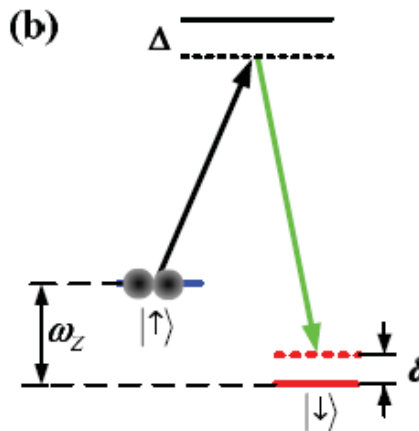
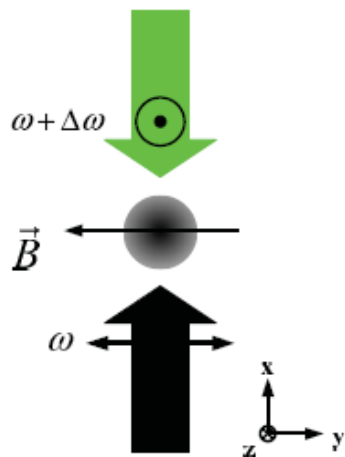
RF spectroscopy of strongly interacting ultracold Fermi gas



$$E_b = \frac{\hbar^2}{m(a - r_0)^2}$$

$$r_0 \approx 60a_0$$

Momentum-resolved **Raman spectroscopy** of non-interacting ultracold Fermi gas



T.-L. Dao, I. Carusotto,
and A. Georges, Phys.
Rev. A 80, 023627 (2009)

$$\hbar \Delta \omega = E_Z + \epsilon^{\text{initial}}(\mathbf{k}) - \epsilon^{\text{final}}(\mathbf{k} + \mathbf{q}_r)$$

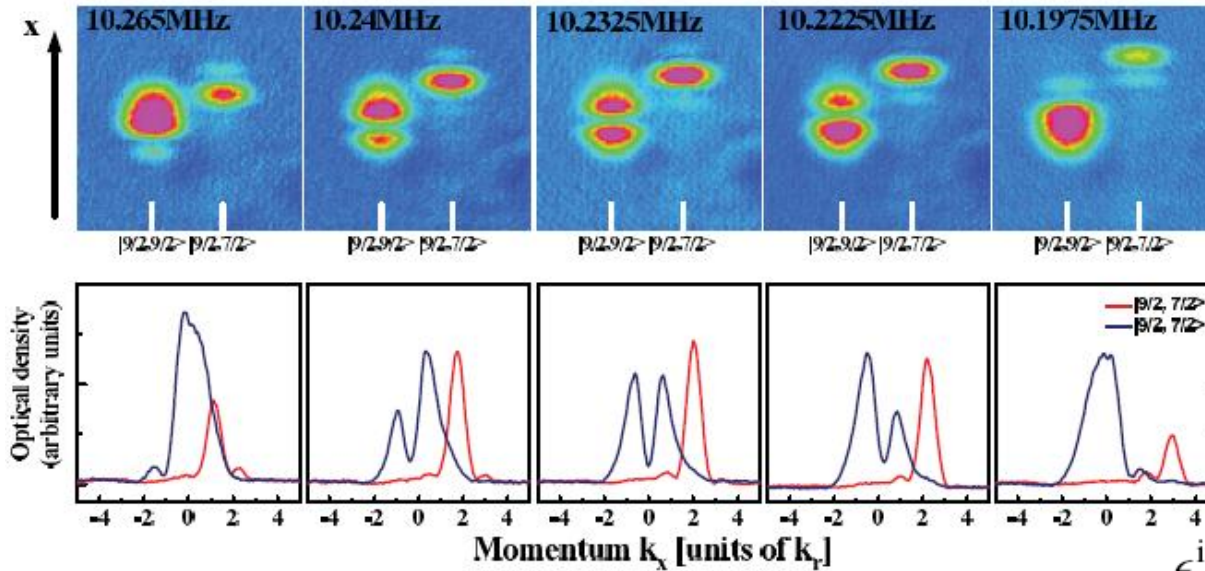
$$|\mathbf{q}_r| = 2k_r \sin(\theta/2)$$

E_Z : the energy split of the two Zeeman states

$\epsilon^{\text{initial}}(\mathbf{k})$: energy-momentum dispersion of the initial state

$\epsilon^{\text{final}}(\mathbf{k} + \mathbf{q}_r)$: energy-momentum dispersion of the final state

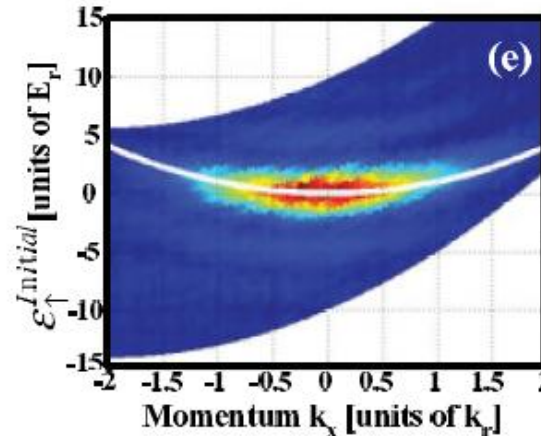
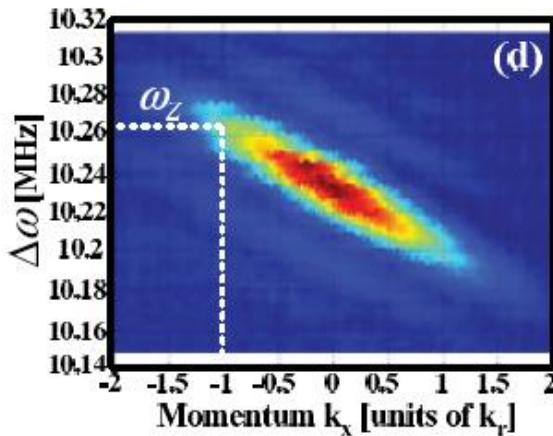
Momentum-resolved Raman spectroscopy of non-interacting ultracold Fermi gas



First prepare fermion in $|9/2\rangle$, and then turn on Raman coupling pulse.

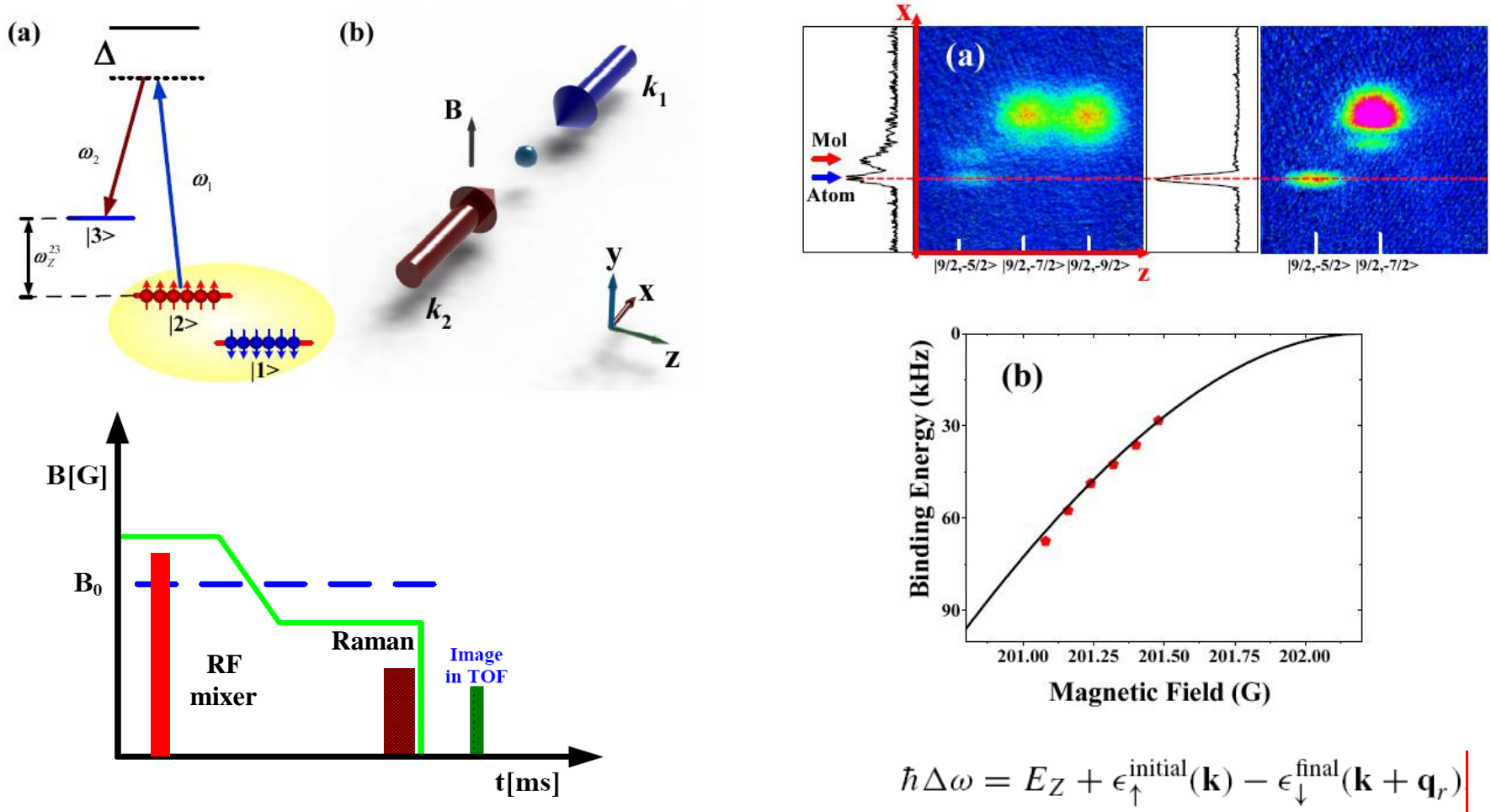
When:

$$\epsilon^{\text{initial}}(\mathbf{k}) = \epsilon^{\text{final}}(\mathbf{k}) = \hbar^2 \mathbf{k}^2 / 2m$$



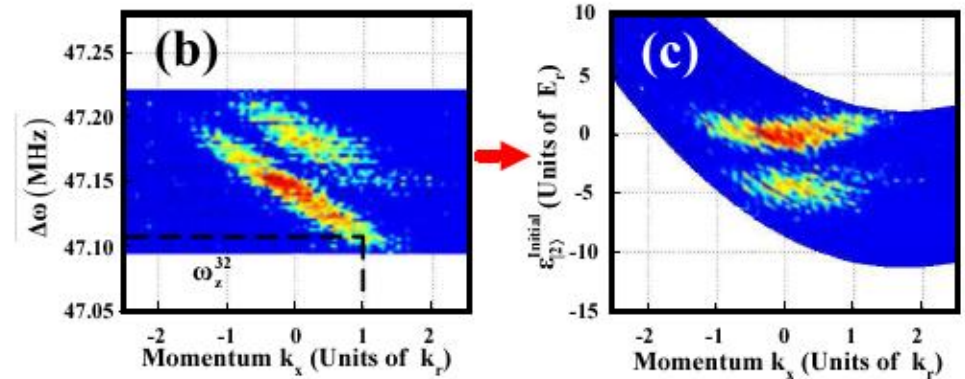
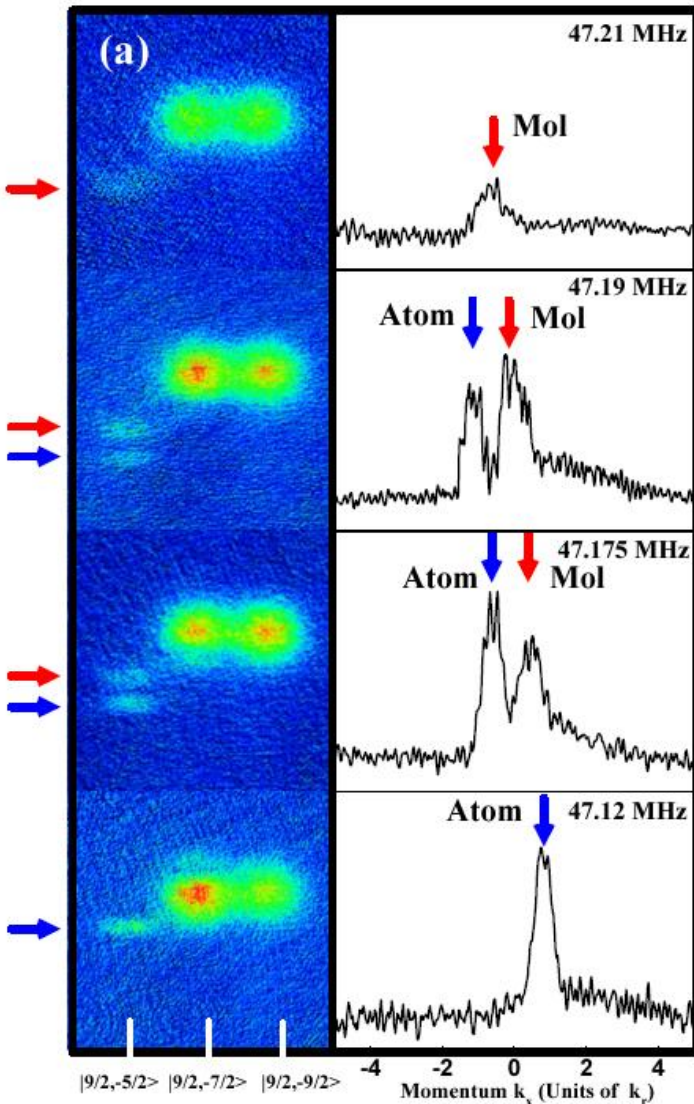
$$\hbar \Delta \omega = E_Z - \frac{\hbar^2 \mathbf{q}_r^2}{2m} - \frac{\hbar^2 \mathbf{q}_r \cdot \mathbf{k}}{m}$$

Momentum-resolved Raman spectroscopy of bound molecules in strongly interacting ultracold Fermi gas



First prepare fermion mixture in $-9/2$ and $-7/2$, and ramp the magnetic field to generate Feshbach molecules, then turn on Raman coupling pulse with Gaussian envelop.

Momentum-resolved Raman spectroscopy of bound molecules in strongly interacting ultracold Fermi gas



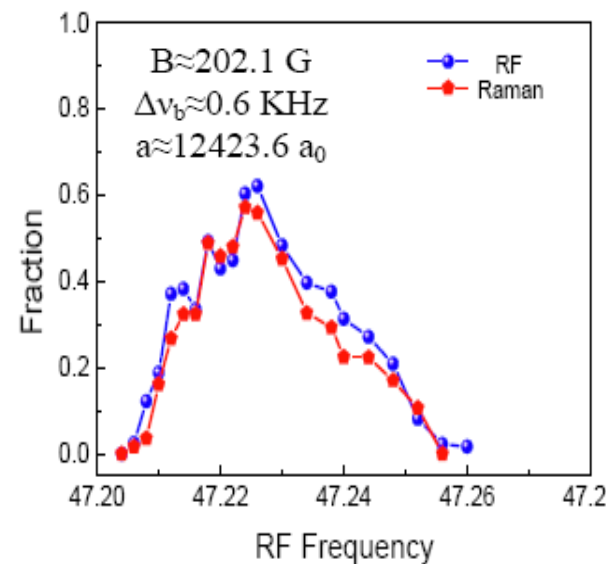
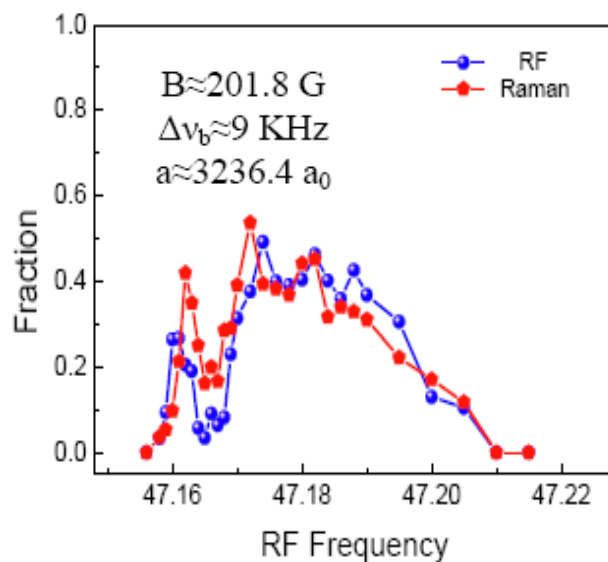
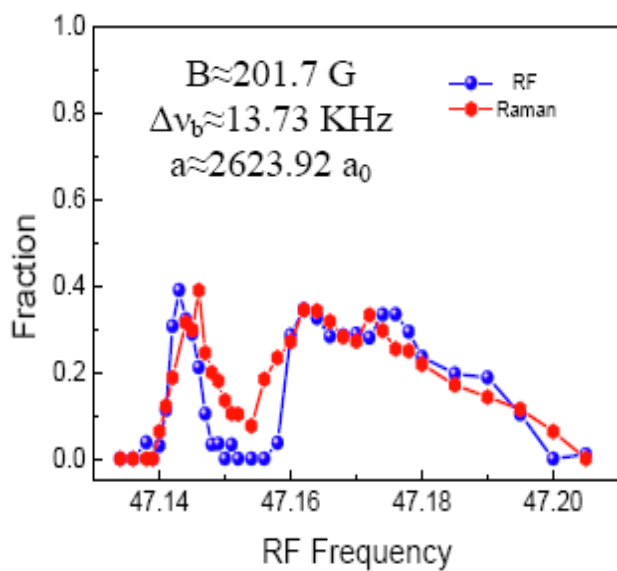
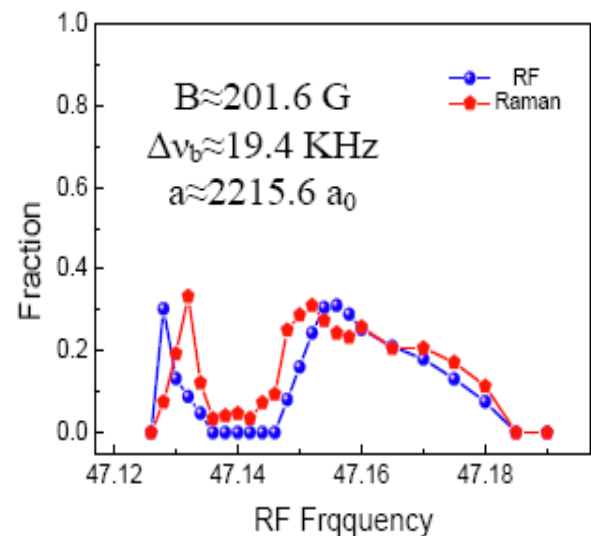
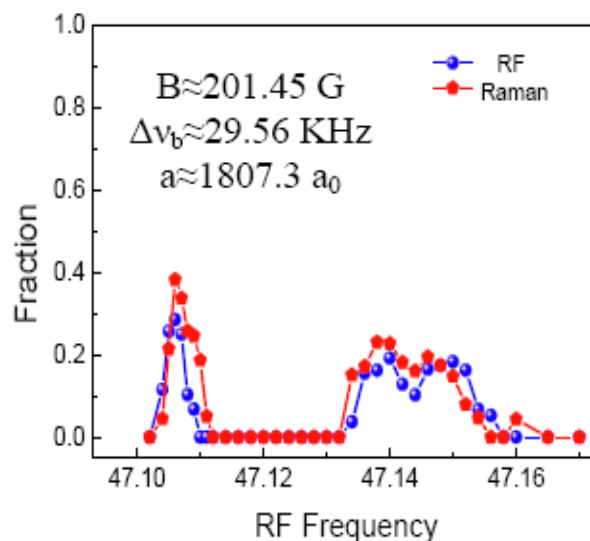
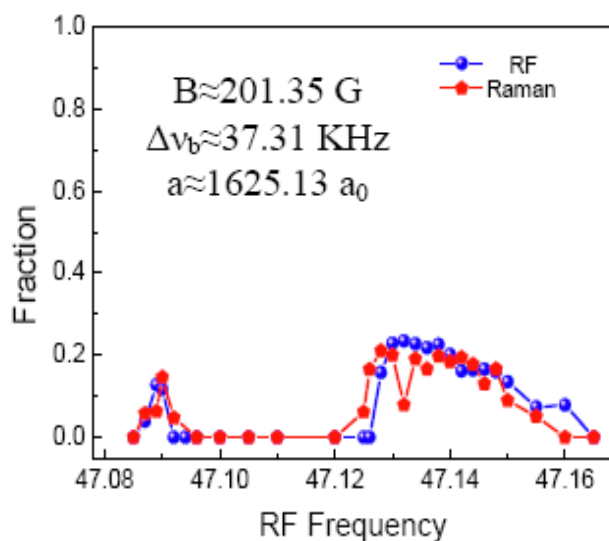
$$\epsilon^{\text{initial}}(\mathbf{k}) = \hbar\Delta\omega - E_Z + \epsilon^{\text{final}}(\mathbf{k} + \mathbf{q}_r)$$

E_Z : the energy split of the two Zeeman states

$\epsilon^{\text{initial}}(\mathbf{k})$: dispersion of the initial state

$\epsilon^{\text{final}}(\mathbf{k} + \mathbf{q}_r)$: dispersion of the final state

Spin-orbit coupling Feshbach molecules



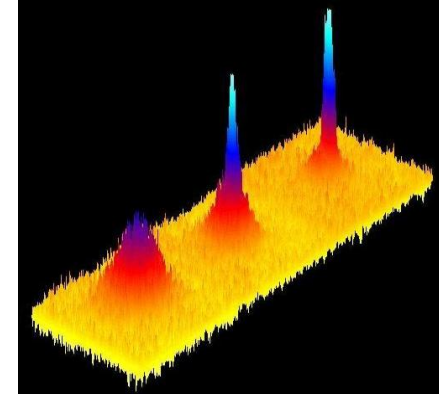


Ultracold⁸⁷Rb-⁴⁰K Bose-Fermi mixture gases

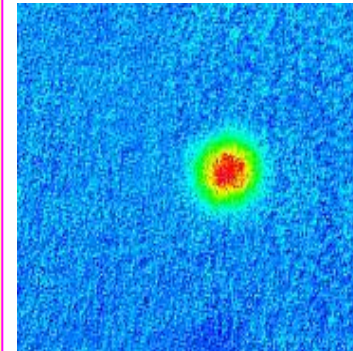
2004. 9 Began to establish

2007.7.7 pm7:00, achieve
⁸⁷Rb BEC

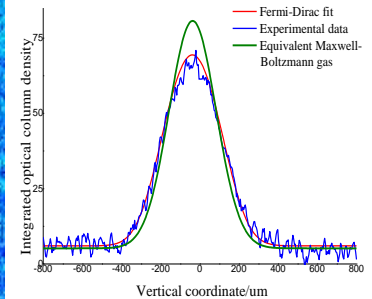
2007.8.30 pm11:00, using
sympathetic cooling
technology to achieve
quantum degenerate of ⁴⁰K
Fermi gas



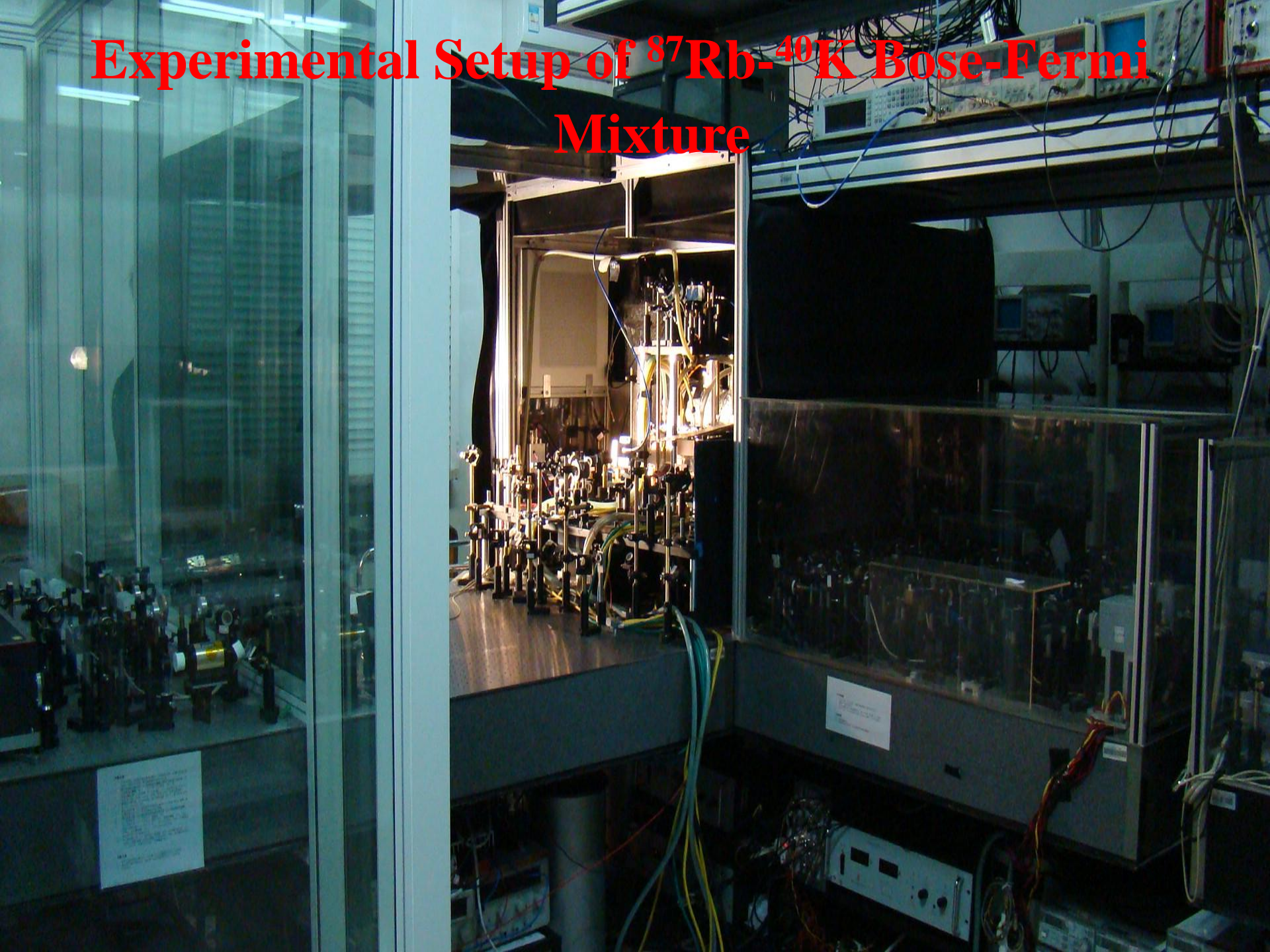
⁸⁷Rb BEC



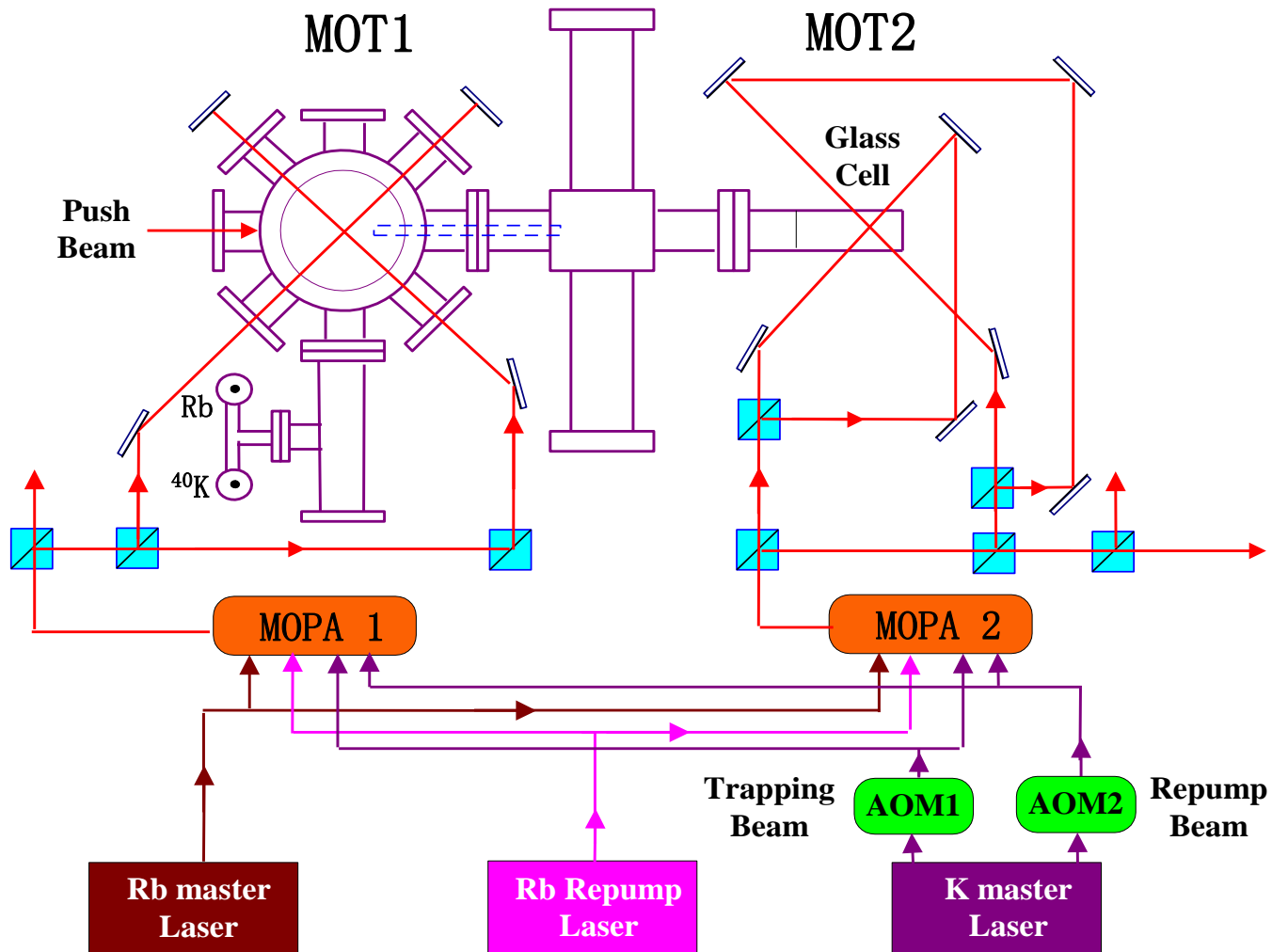
⁴⁰K DFG



Experimental Setup of ^{87}Rb - ^{40}K Bose-Fermi Mixture



Schematic of experimental setup



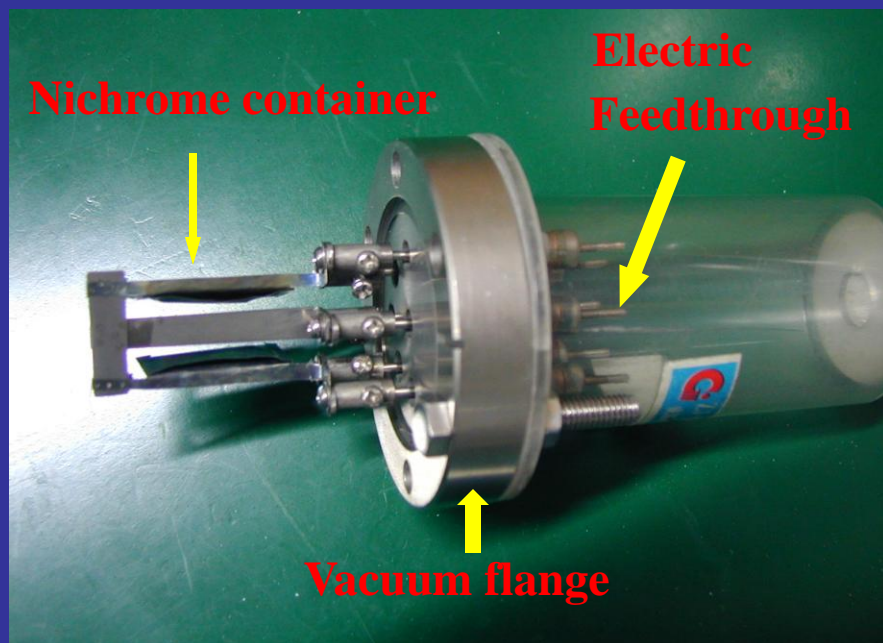
Potassium 40 Dispenser



B. DeMarco, H. Rohner, and D. S. Jin, Rev. Sci. Instrum. 70, 1967 (1999).

0.0012% ^{40}K naturally

Enriched 6% ^{40}K

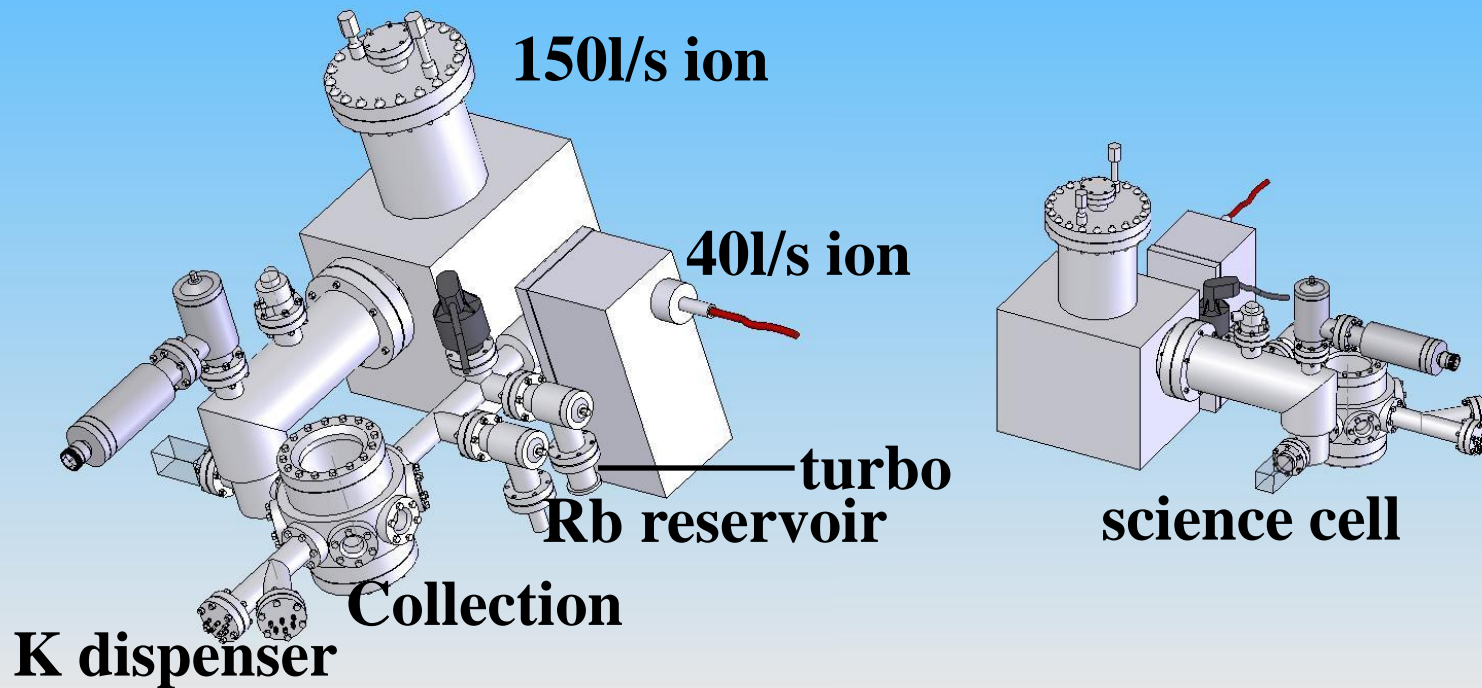


Dispenser



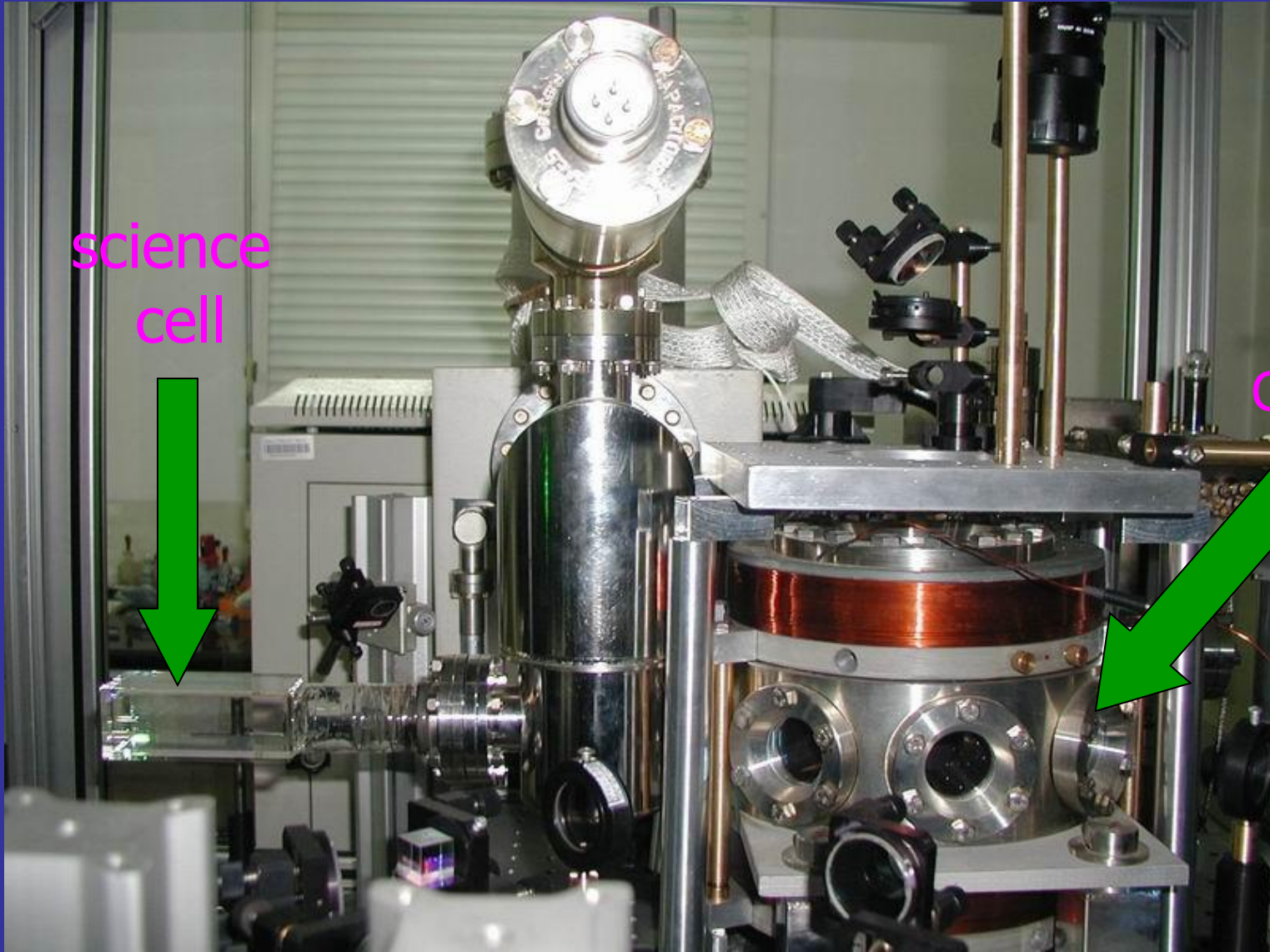
Glove Box

Vacuum system



First Chamber: $1.2 \cdot 10^{-7} \text{Pa}$

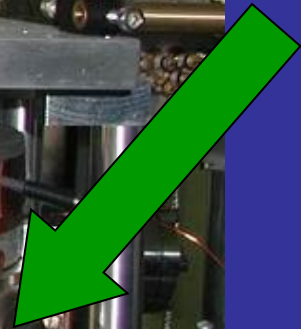
Second Chamber: $2.9 \cdot 10^{-9} \text{Pa}$



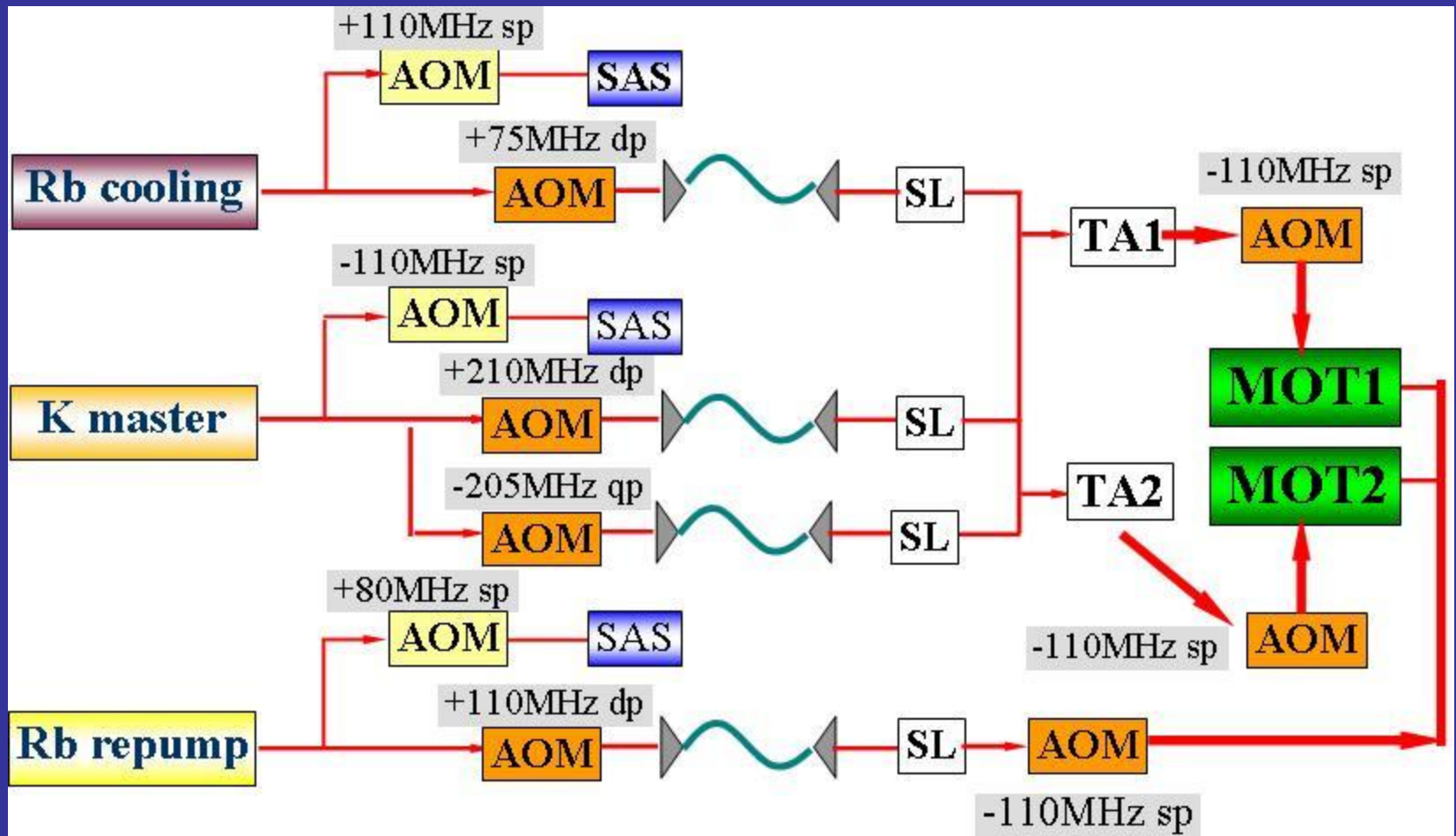
science
cell



Collection
cell



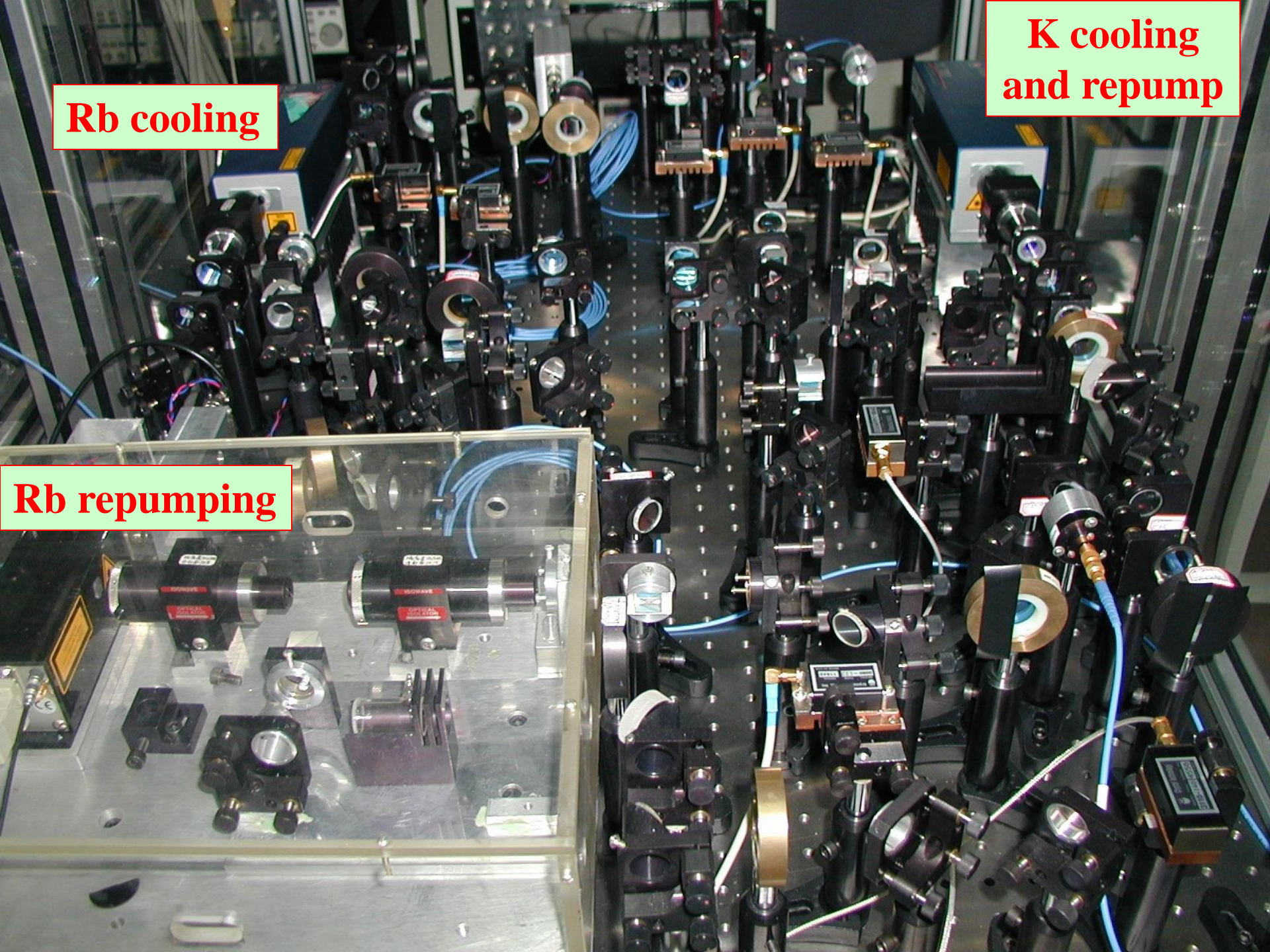
Laser System

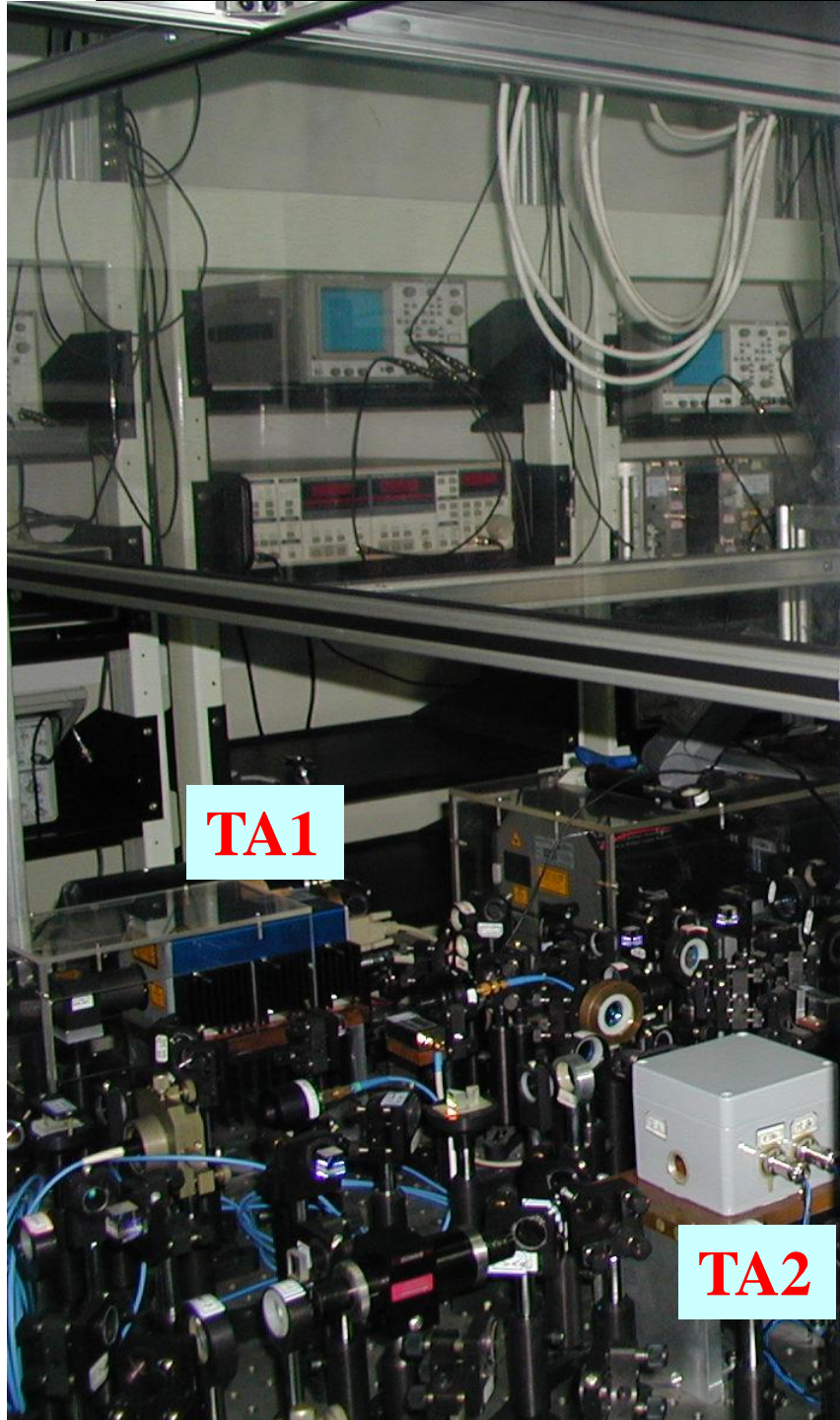


Rb cooling

**K cooling
and repump**

Rb repumping





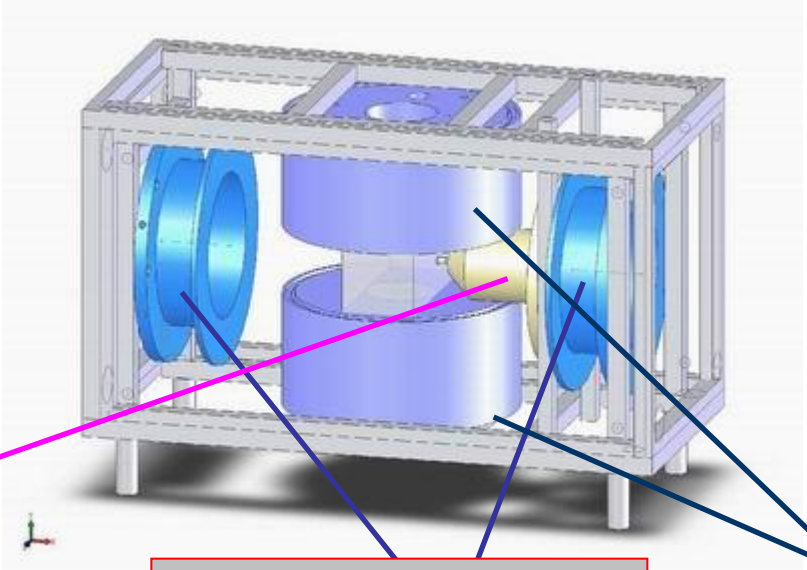
TA1

TA2



Injection locking slave laser

Magnetic trap: Quadrupole-Ioffe configuration (Quic) trap



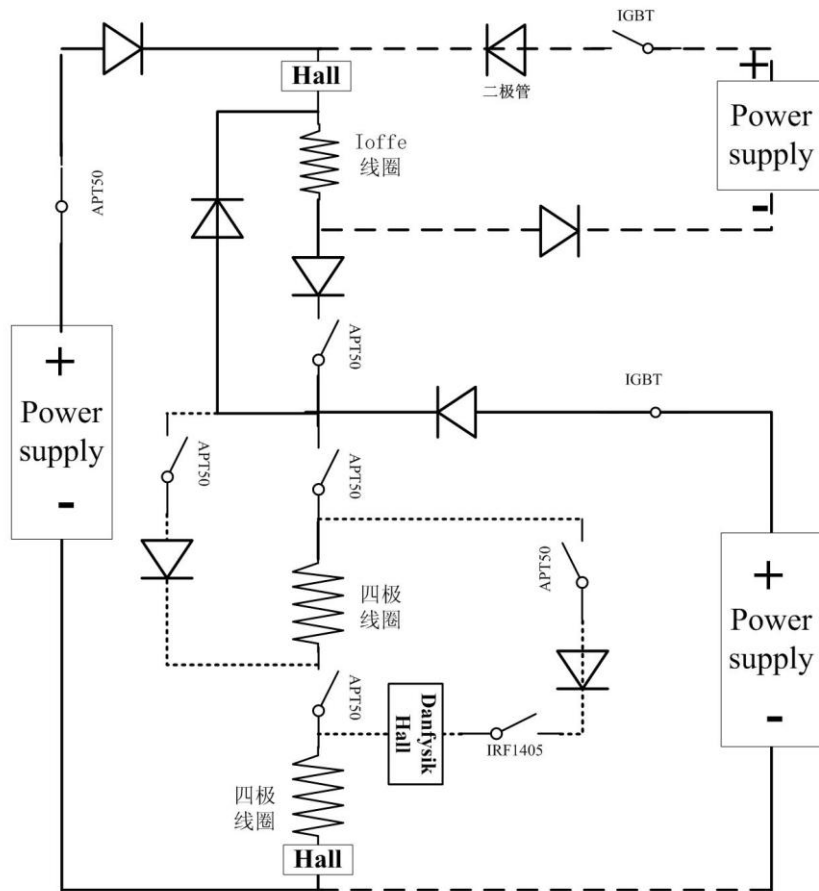
Ioffe coil

Bias coils

Quadrupole coils



Control magnetic field current



- MOT2
- QUIC trap
- Feshbach resonance



Experimental Sequence

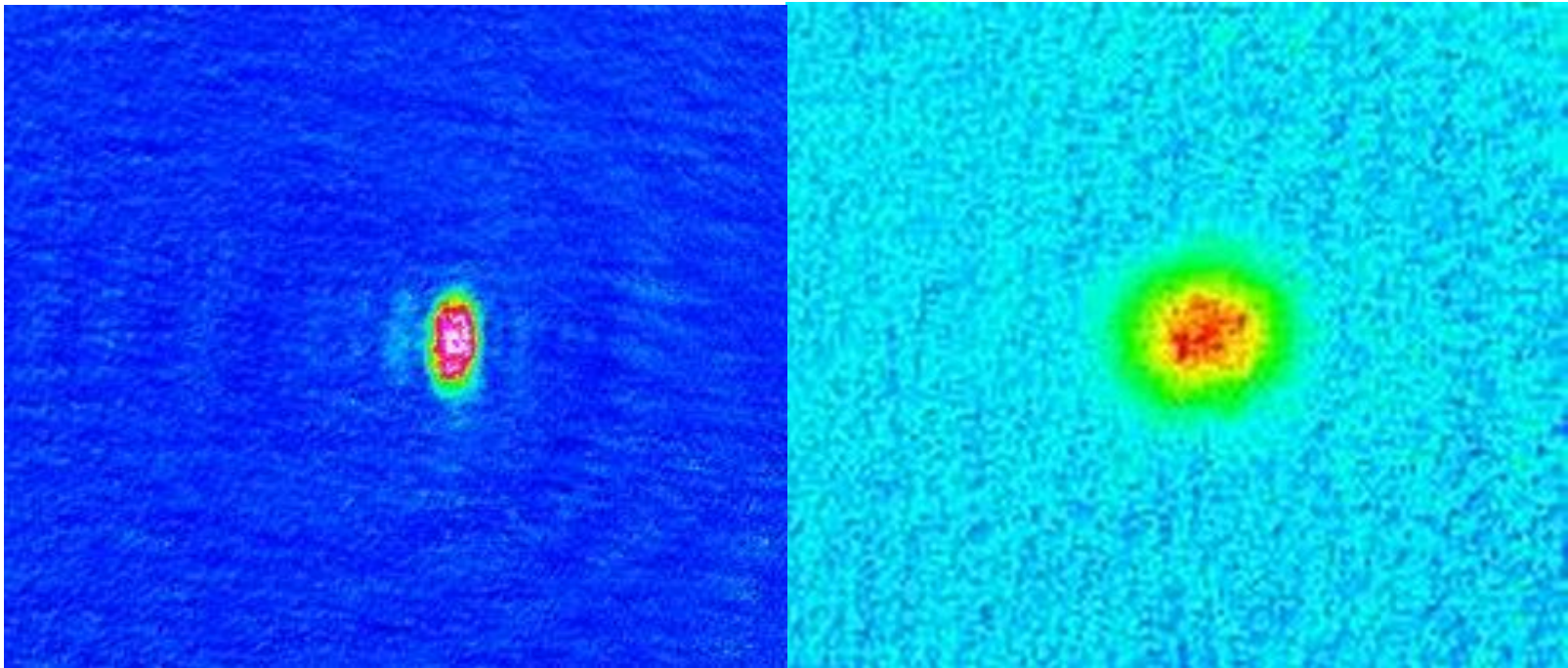
- **Transfer MOT1 to MOT2**
- **Optical molasses and Optical pump**
- **Loading atoms into QUIC
(Quadrupole Ioffe Configuration)**
- **RF evaporation**
- **Time of flight and absorption image**



Generation of quantum degenerate gases

^{87}Rb BEC

^{40}K DFG

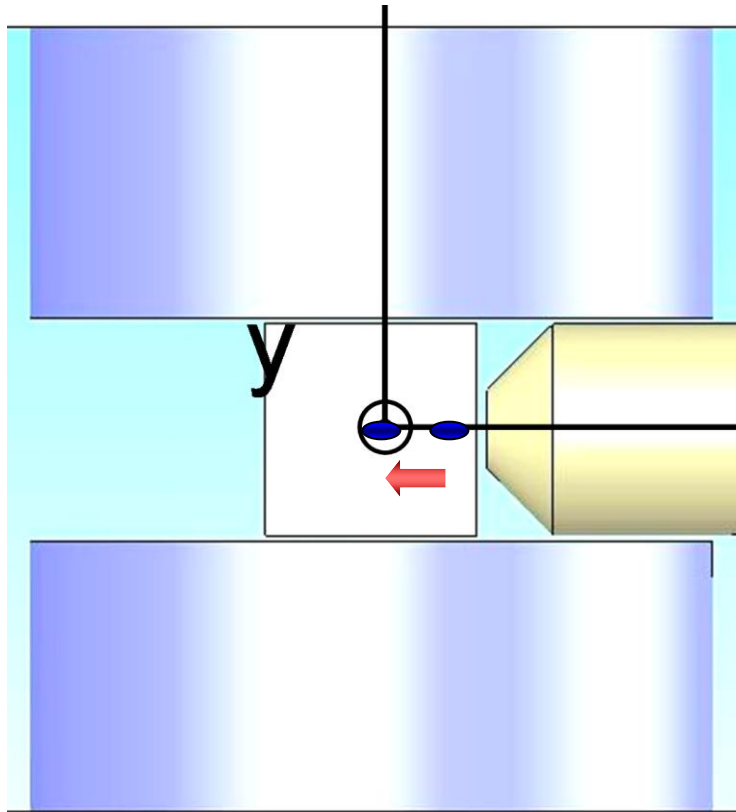


$N=1.8 \cdot 10^5$
 $t_f=25\text{ms}$

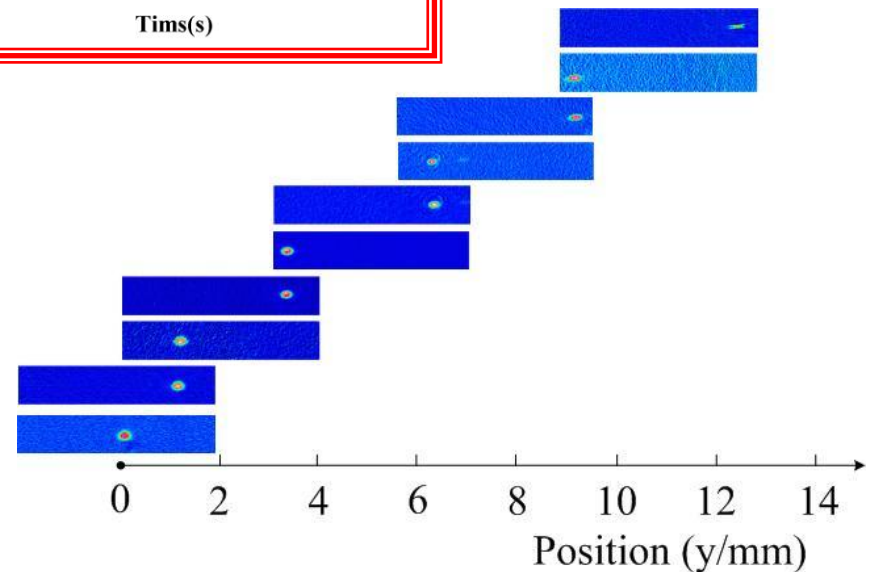
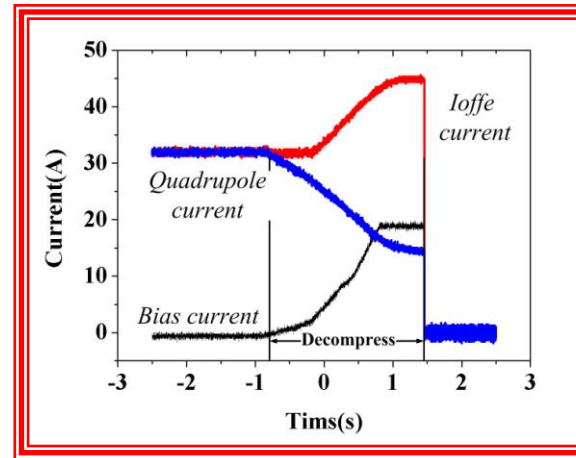
$T/T_F=0.56$
 $N=5.5 \cdot 10^5$
 $t_f=12\text{ms}$

Dezhi Xiong, Haixia Chen, Pengjun Wang, Xudong Yu, Feng Gao, Jing Zhang, *Chin. Phys. Lett.* 25, 843 (2008).

Transfer ultracold atoms from QUIC to the center of glass cell

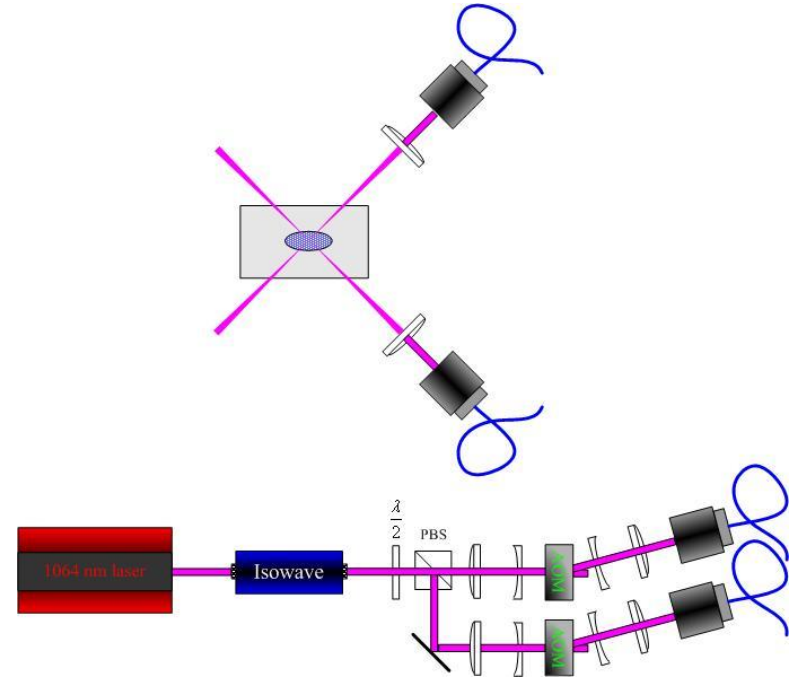
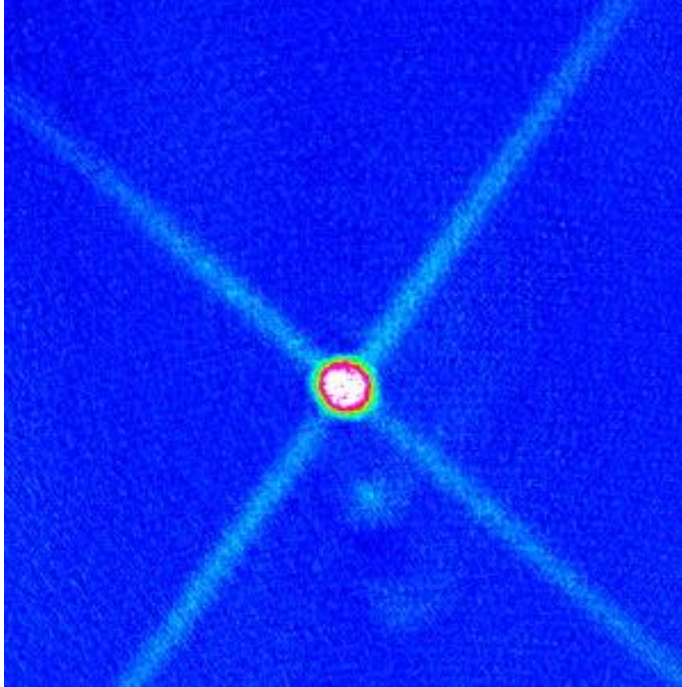


Distance: 11.8 mm



Dezhi Xiong, Haixia Chen, Pengjun Wang, Zhengkun Fu, Jing Zhang, *Opt. Express* 18, 1649 (2010)

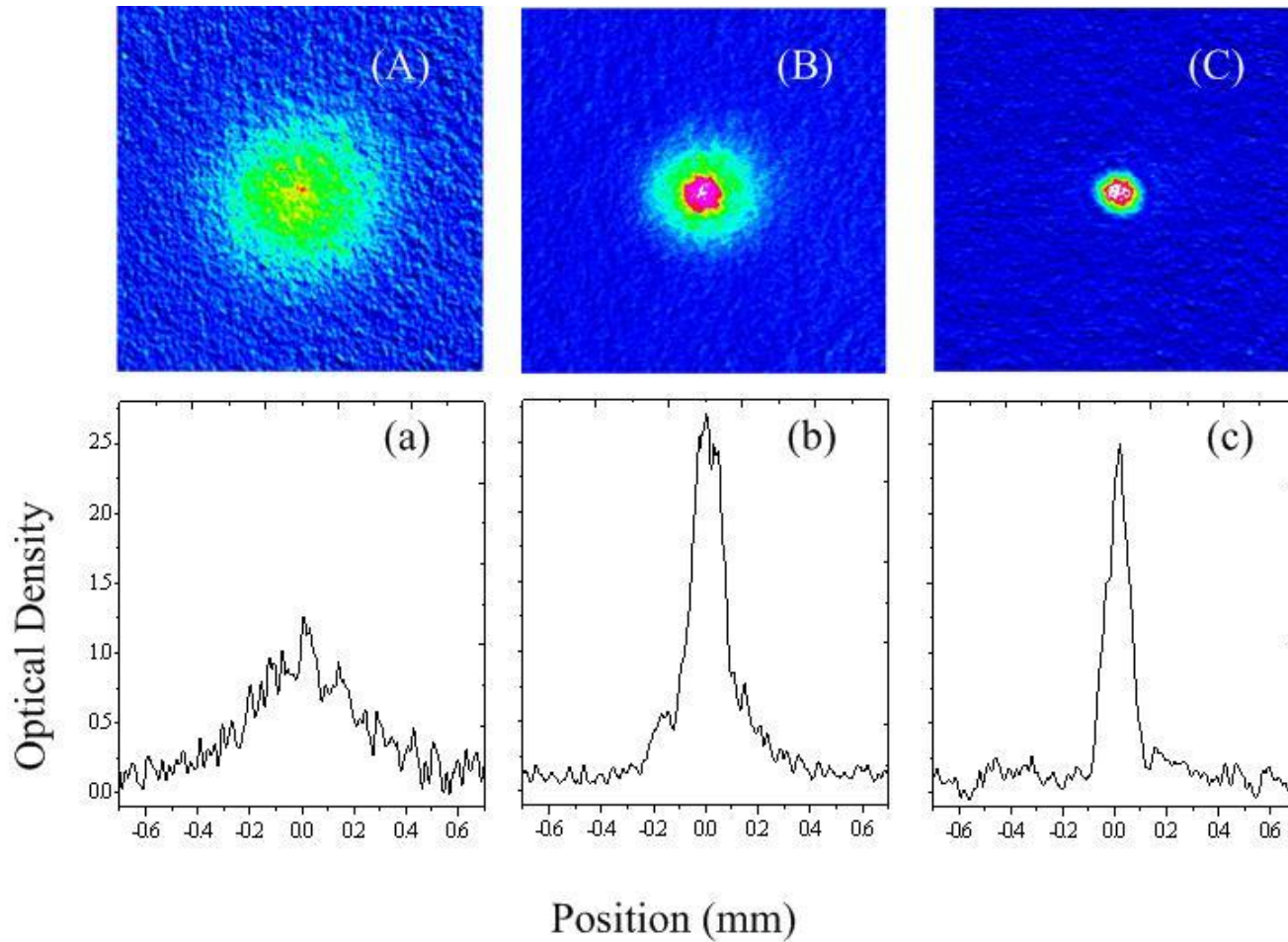
Loading of a crossed dipole trap



$P1=380\text{mw}$ $P2=650\text{mw}$

$W1=25\mu\text{m}$ $W2=33\mu\text{m}$

1064 nm single-frequency laser



$N=1.5 \times 10^5$

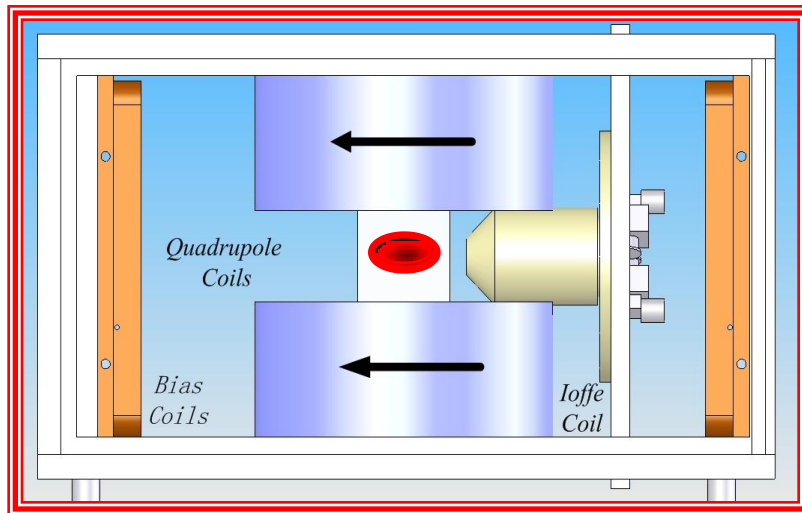
30 ms free expansion time.

Feshbach resonance

Beginning: Evaporation in optical trap

Number: ^{87}Rb & ^{40}K --- $7-8(10) \cdot 10^5$

Temperature: $1-2 \mu\text{k}$



❖ Feshbach resonance

Key point:

Stability of magnetic field

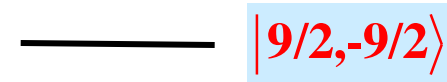
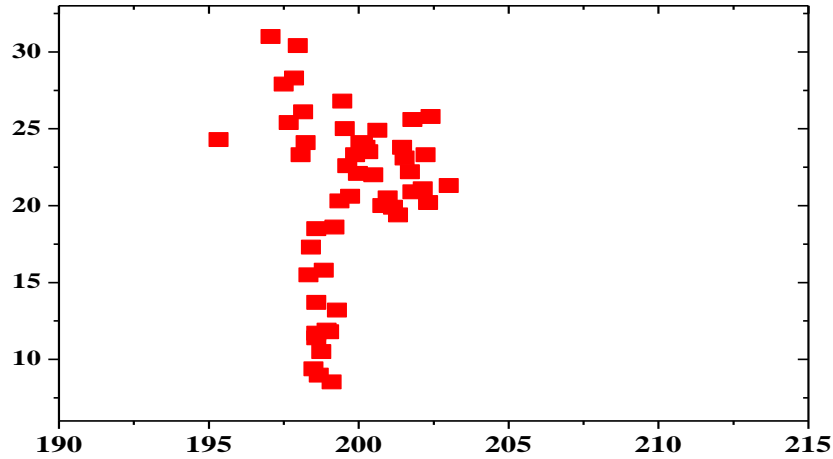
1. Prepare of spin state

2. Ramp magnetic field

3. Measure the atomic losses and heating

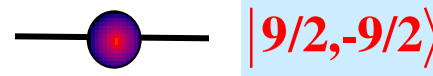
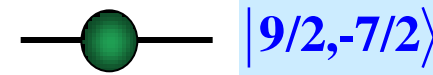
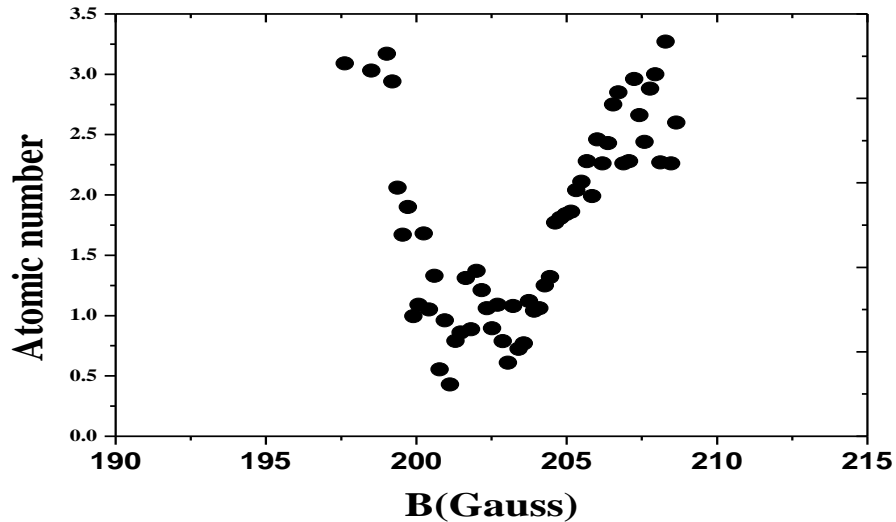
4. Position and width of Feshbach resonance

Observed Feshbach resonances in 40K



P-wave $B_0 = 198.8$ G

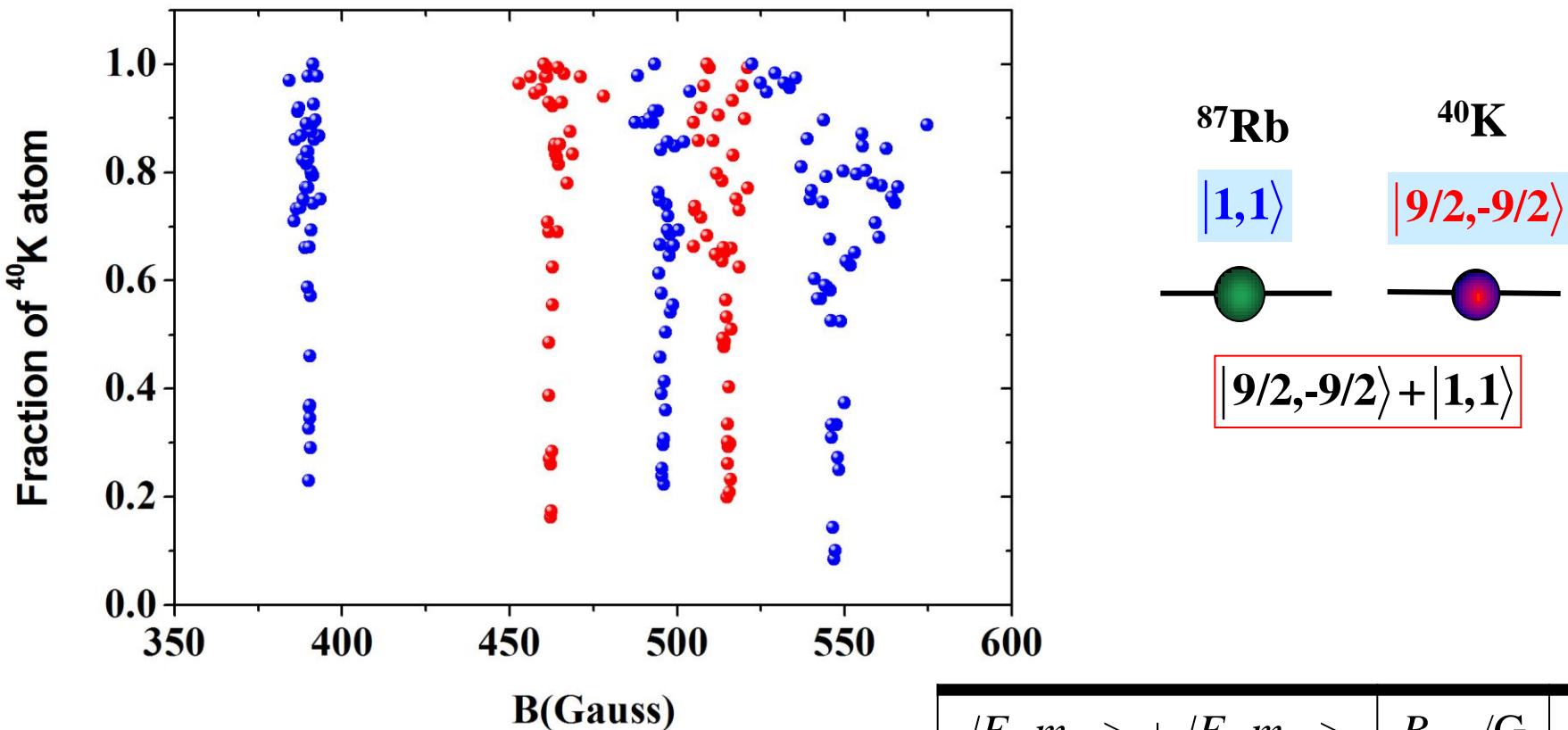
$|9/2, -7/2\rangle$



S-wave $B_0 = 202.1$ G

$|9/2, -9/2\rangle + |9/2, -7/2\rangle$

Observed Feshbach resonances in $^{40}\text{K}+^{87}\text{Rb}$

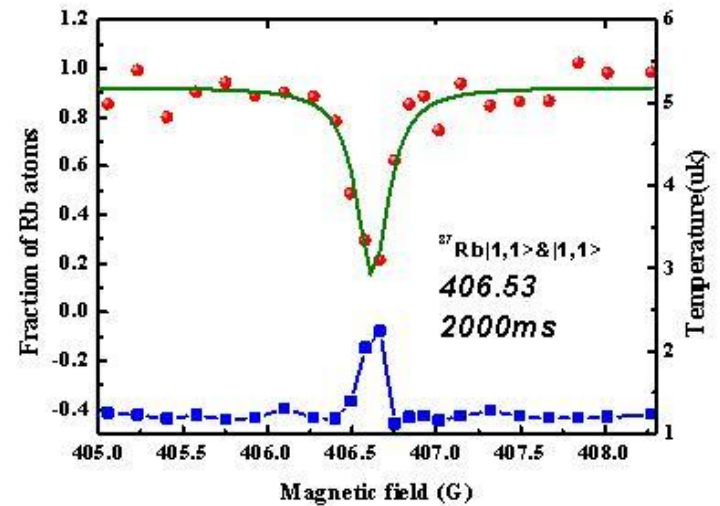
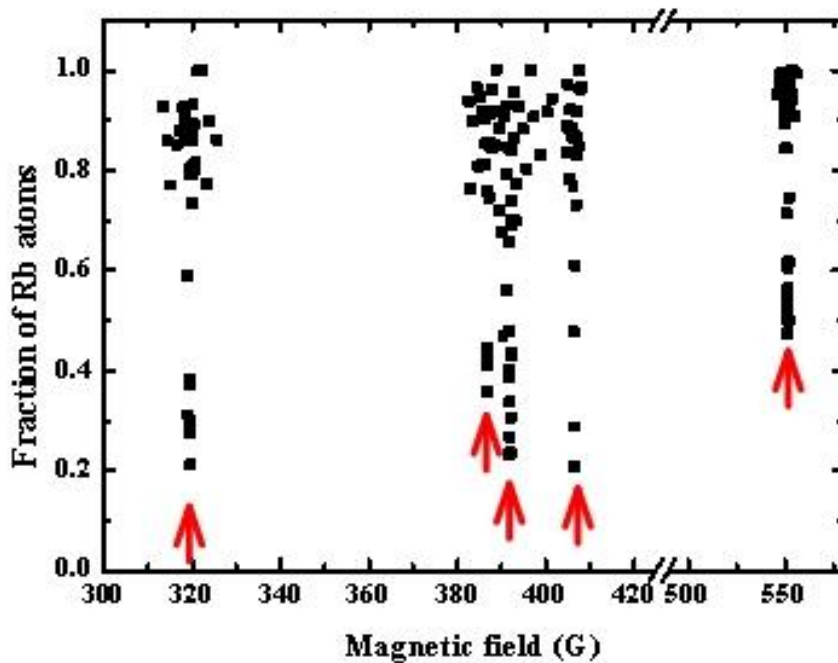


$ F_1, m_{F1}\rangle + F_2, m_{F2}\rangle$	B_{expt}/G	L
$ 1,1\rangle + 9/2,-9/2\rangle$	462.45	0
	495.71	0
	515.75	1
	546.89	0

❖ Feshbach resonance

Feshbach resonance ^{87}Rb $|1,1\rangle$ & $|1,1\rangle$

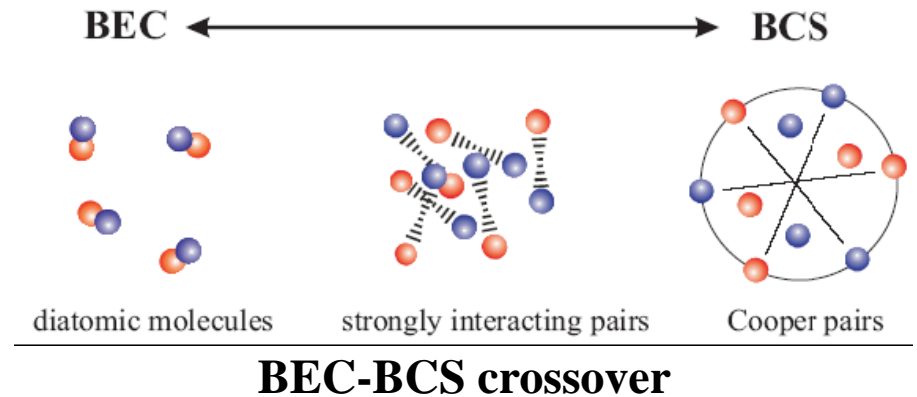
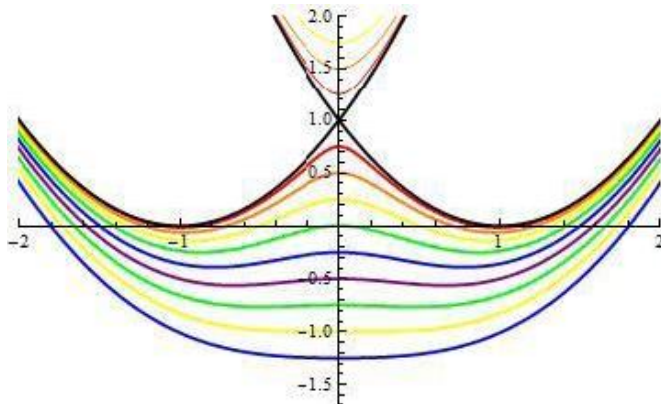
$ F_1, m_{F1}\rangle + F_2, m_{F2}\rangle$	B_{expt}/G	l
$ 1,1\rangle + 1,1\rangle$	319.3	2
	387.25	2
	406.54	0
	551.47	2



Future Plan

BEC-BCS crossover with spin-orbit coupling

Method: RF and Raman spectrum



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Shijie Chai

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Zengming Meng

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Zengqiang Yu

Jiao Miao

Swinburne University of Technology

Hui Hu, Xiaji Liu

Rice University

Han Pu





Thank you!





Thank you!