

Phase transitions in a strongly interacting Bose gas

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Outline

1. Introduction

- (1) Traditional theory of dilute Bose gases
- (2) Results about strongly interacting Bose gases

2. Instability in the atomic condensation

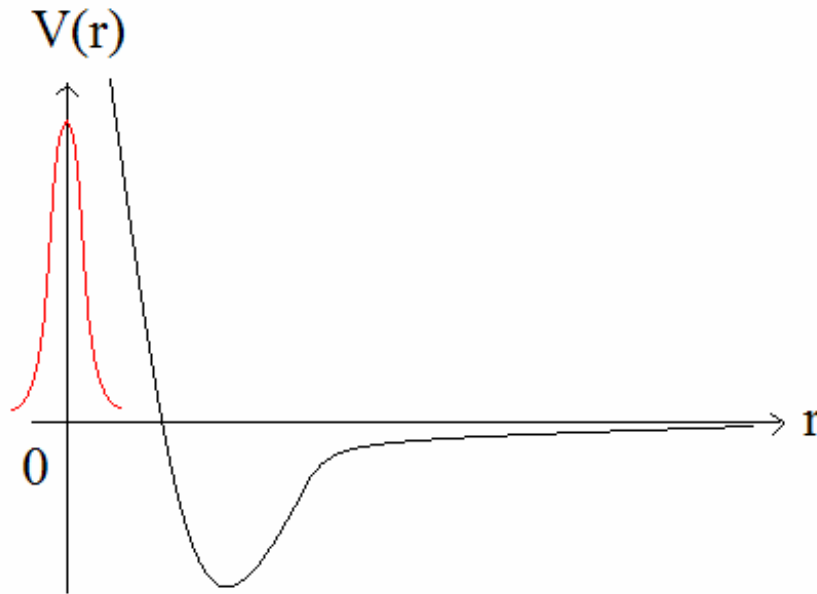
3. The molecular condensation state

4. Coherent mixture of atoms and molecules

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1. Introduction

(1) Tradition theory of dilute Bose gases



S-wave phase shift

$$\delta_0 \rightarrow -ka$$

Renormalization of the hard-core potential

$$V(\mathbf{r}) = 4\pi \frac{\hbar^2}{m} a \delta^3(r)$$

(Fermi, Lee & Yang & Huang, Galitskii, Beliaev)

Beliaev's Theory (Homogeneous, T=0):

- The effective potential and the scattering amplitude satisfy similar integral equations.

$$V \begin{array}{|c|} \hline \diagup \\ \hline \end{array} = \text{---} + \begin{array}{|c|} \hline \text{---} \\ \hline \end{array} + \begin{array}{|c|} \hline \text{---} \\ \hline \end{array} + \dots,$$

- Perturbation in gas parameter $\sqrt{8\pi n_0 a^3}$.

- Beliaev damping of the phonon mode $\Gamma_k = \frac{3\hbar^2 k^5}{640\pi m n_0}$.
(One loop)

Hugenholtz and Pines Theorem:

The quasi-particle energy is gapless. (Valid in all orders)

Bogoliubov's Theory:

Hamiltonian $H = -\psi^\dagger \frac{\hbar^2 \nabla^2}{2m} \psi + \frac{g}{2} \psi^\dagger \psi^\dagger \psi \psi$

Condensation $\langle \psi \rangle = \psi_0$, and fluctuation $\delta\psi \equiv \psi - \psi_0$

Quadratic Hamiltonian

$$H_2 = \sum_{\mathbf{k}} (\varepsilon_{\mathbf{k}} + gn_0) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{gn_0}{2} \sum_{\mathbf{k}} (\psi_{-\mathbf{k}}^\dagger \psi_{\mathbf{k}}^\dagger + \psi_{\mathbf{k}} \psi_{-\mathbf{k}})$$

Applying Bogoliubov Transformation $b_{\mathbf{k}} = u_{\mathbf{k}} \psi_{\mathbf{k}} - v_{\mathbf{k}} \psi_{-\mathbf{k}}^\dagger$

$$H_2 = C + \sum_{\mathbf{k}} E_{\mathbf{k}} b_{\mathbf{k}}^\dagger b_{\mathbf{k}}$$

Quasi-particle energy $E_{\mathbf{k}} = \sqrt{\varepsilon_{\mathbf{k}} (\varepsilon_{\mathbf{k}} + 2gn_0)}$

Quantum depletion $n - n_0 = \frac{1}{V} \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} \rangle = \frac{8}{3} \sqrt{\frac{n_0^3 a^3}{\pi}}$

(2) Results about strongly interacting Bose gases

Feshbach resonance

Atom $\mu_a \approx \mu_B$ (triplet)

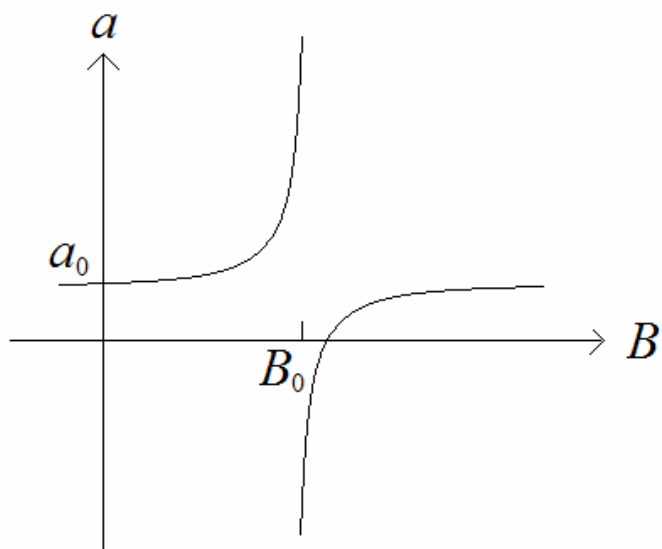
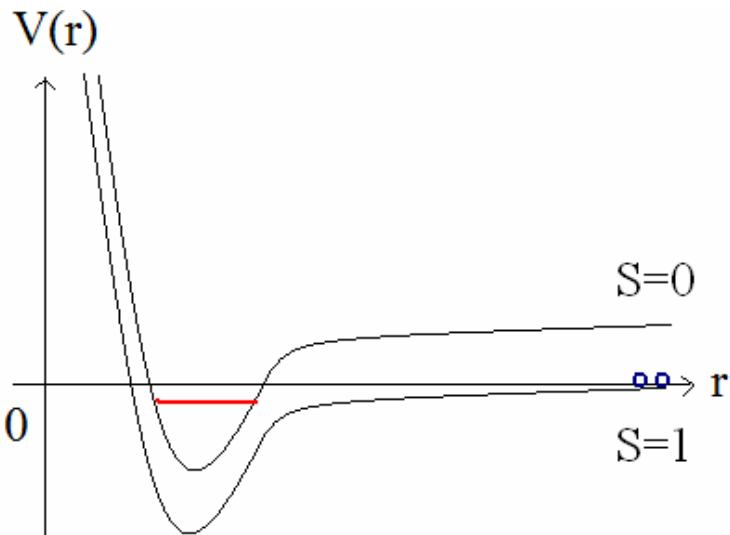
Diatomic molecule $\mu_m \approx 0$ (singlet)

Zeeman-energy difference

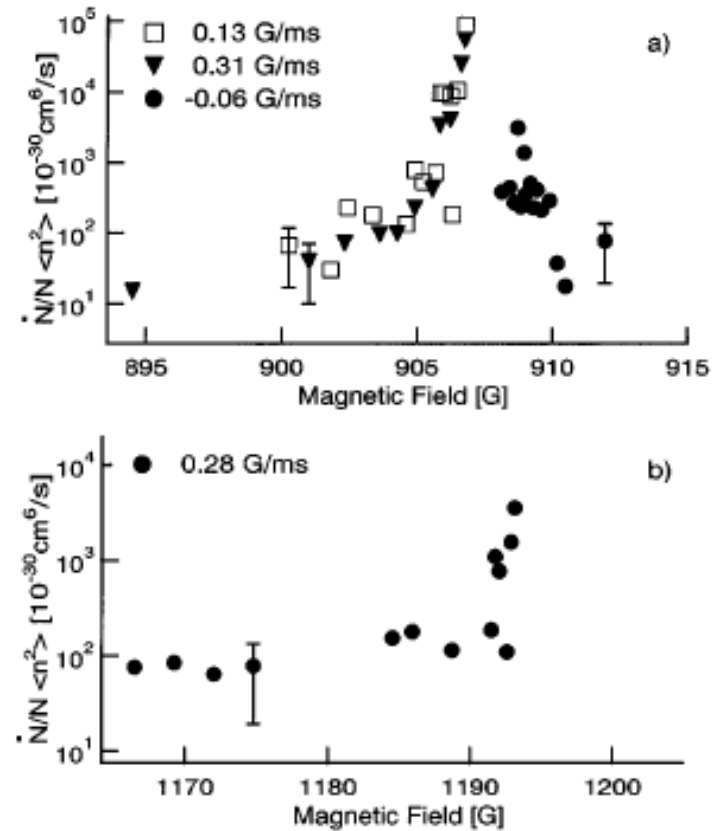
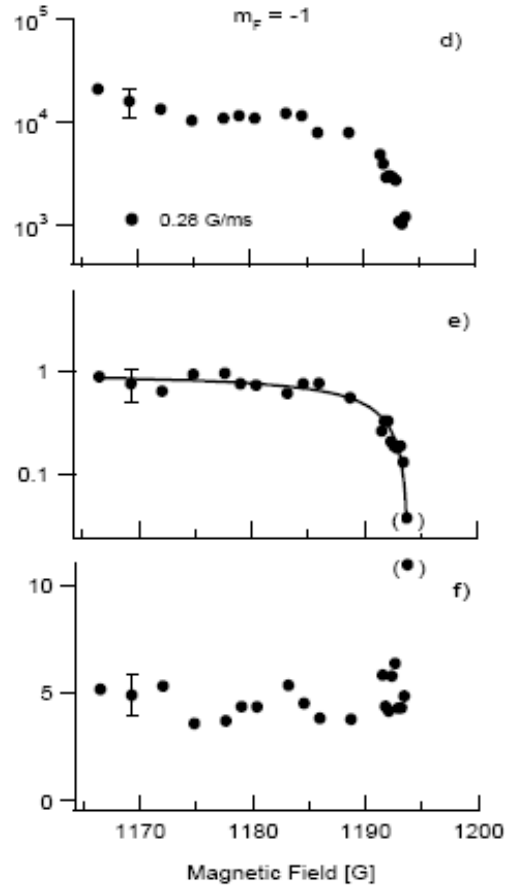
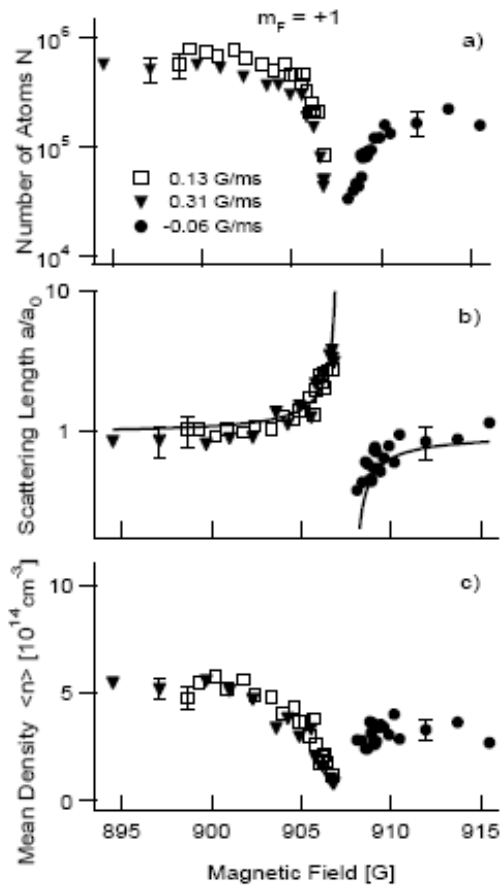
$$\Delta E_Z \approx \mu_B B$$

Scattering length can be tuned by magnetic field.

$$a = a_0 \left(1 - \frac{\Delta}{B - B_0}\right)$$

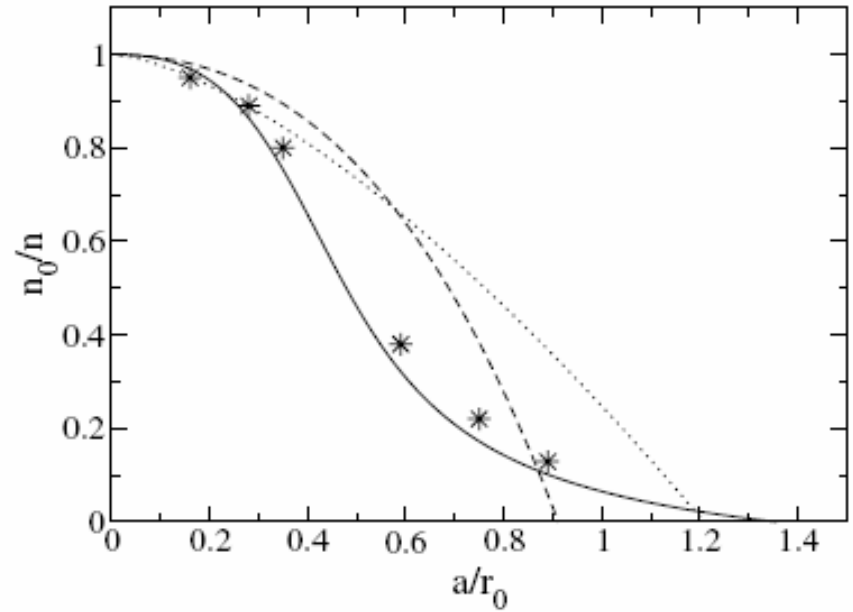
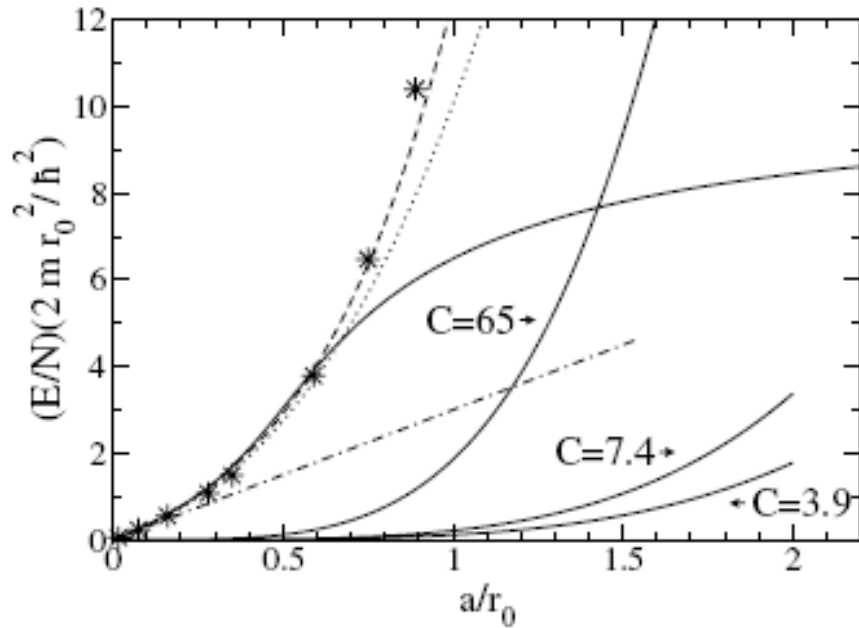


Large atom loss rate near resonance



(Ketterle *et al* 98, 99)

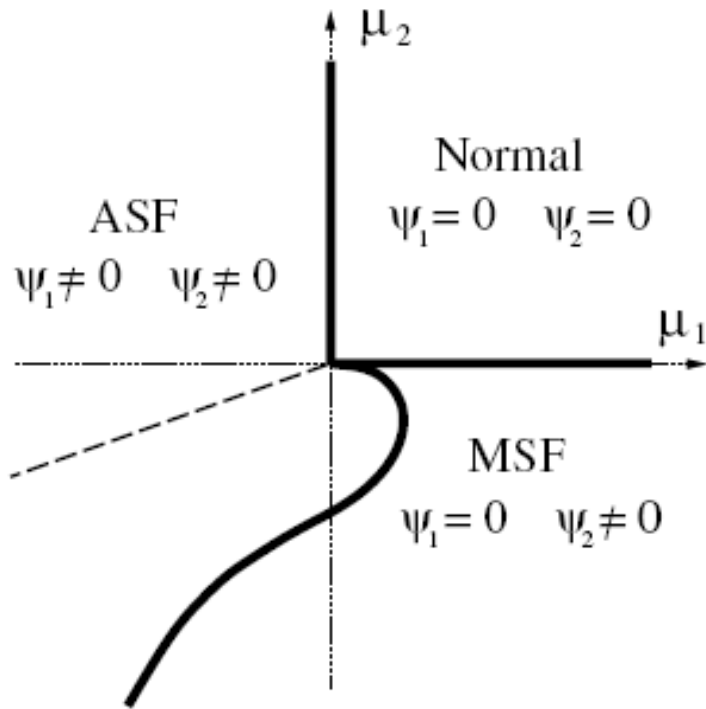
LOCV Result



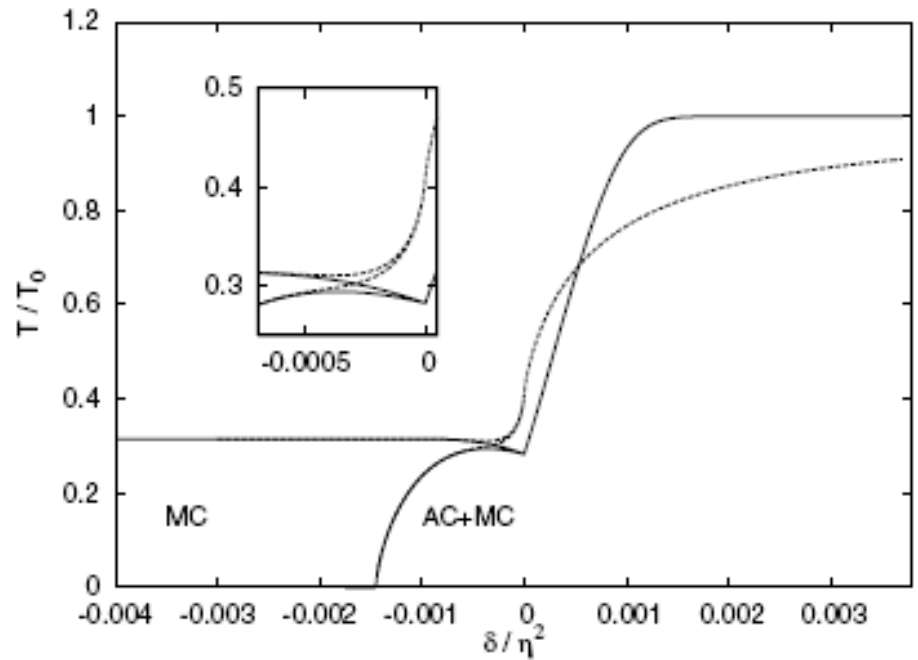
At resonance
$$E/N = \frac{\lambda}{2} = \left(\frac{3}{2}\right)^{2/3} \frac{\hbar^2 (kd)^2}{2mr_0^2} = 13.33 \hbar^2 \frac{n^{2/3}}{m}.$$

(Pethick *et al* 02)

Transition between atomic and molecular condensations



(Radzihovsky *et al* 04)

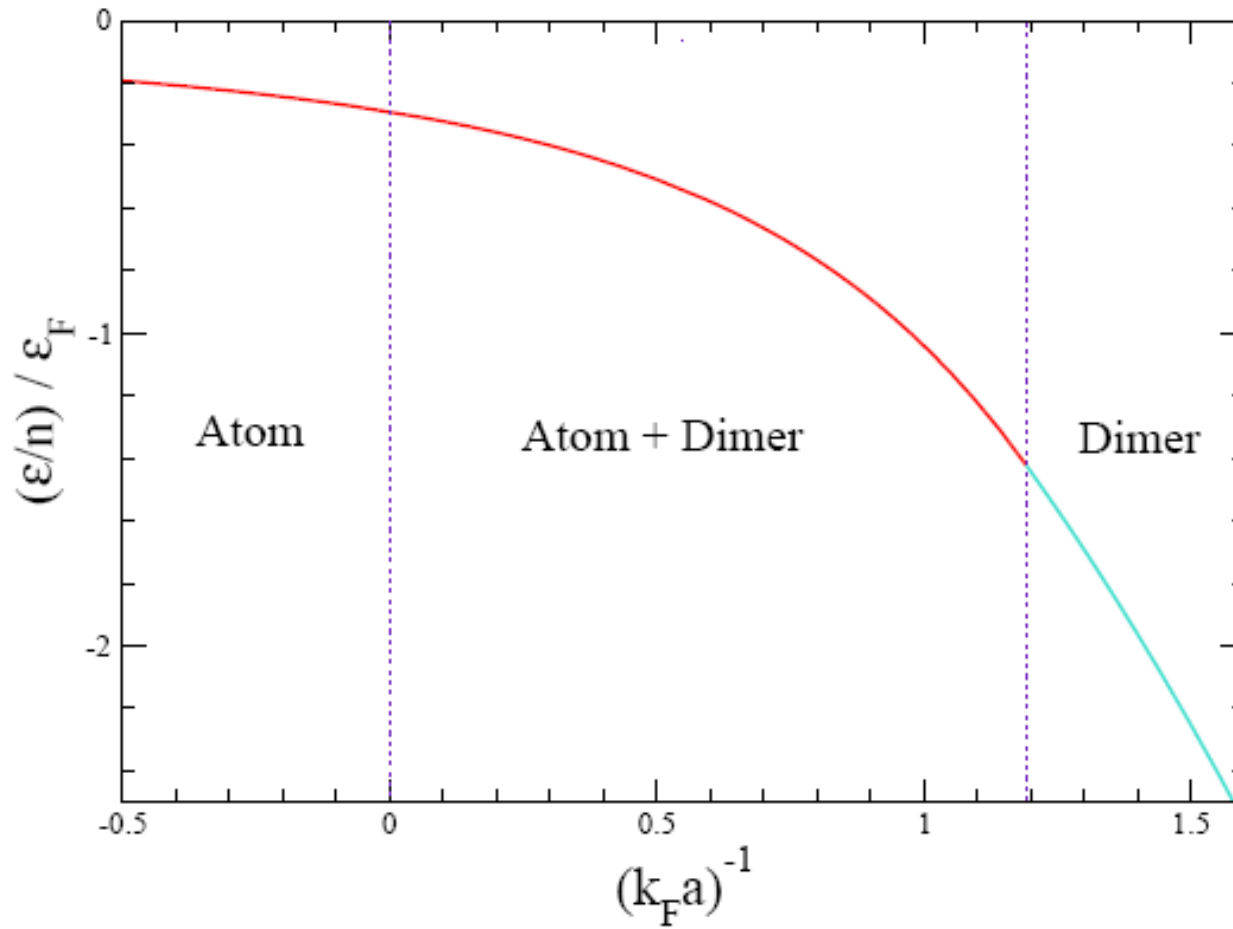


(Stoof *et al* 04)

Negative compressibility of atomic condensation ($a < 0$)

(Mueller *et al* 05)

Braaten's Theory (07)



2. Instability of Atomic condensation ($a > 0$)

Hamiltonian $H = -\psi^\dagger \frac{\hbar^2 \nabla^2}{2m} \psi + \frac{g}{2} \psi^\dagger \psi^\dagger \psi \psi$, $g \equiv 4\pi \frac{\hbar^2 a}{m}$

Condensation $\langle \psi \rangle = \psi_0$, and fluctuation $\delta\psi \equiv \psi - \psi_0$

Quadratic Hamiltonian

$$H_2 = \sum_{\mathbf{k}} (\varepsilon_{\mathbf{k}} + gn_0) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{gn_0}{2} \sum_{\mathbf{k}} (\psi_{-\mathbf{k}}^\dagger \psi_{\mathbf{k}}^\dagger + \psi_{\mathbf{k}} \psi_{-\mathbf{k}})$$

Bogoliubov Transformation $b_{\mathbf{k}} = u_{\mathbf{k}} \psi_{\mathbf{k}} - v_{\mathbf{k}} \psi_{-\mathbf{k}}^\dagger$

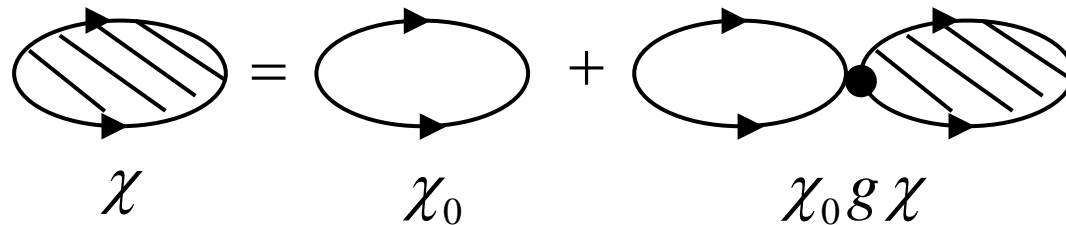
$$H_2 = C + \sum_{\mathbf{k}} E_{\mathbf{k}} b_{\mathbf{k}}^\dagger b_{\mathbf{k}}$$

Quasi-particle energy $E_{\mathbf{k}} = \sqrt{\varepsilon_{\mathbf{k}} (\varepsilon_{\mathbf{k}} + 2gn_0)}$

Quantum depletion $n - n_0 = \frac{1}{V} \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} \rangle = \frac{8}{3} \sqrt{\frac{n_0^3 a^3}{\pi}}$

The spectrum of molecular excitations are given by the poles of two-particle correlation functions.

In RPA approximation $\chi = (1 - g \chi_0)^{-1} \chi_0$

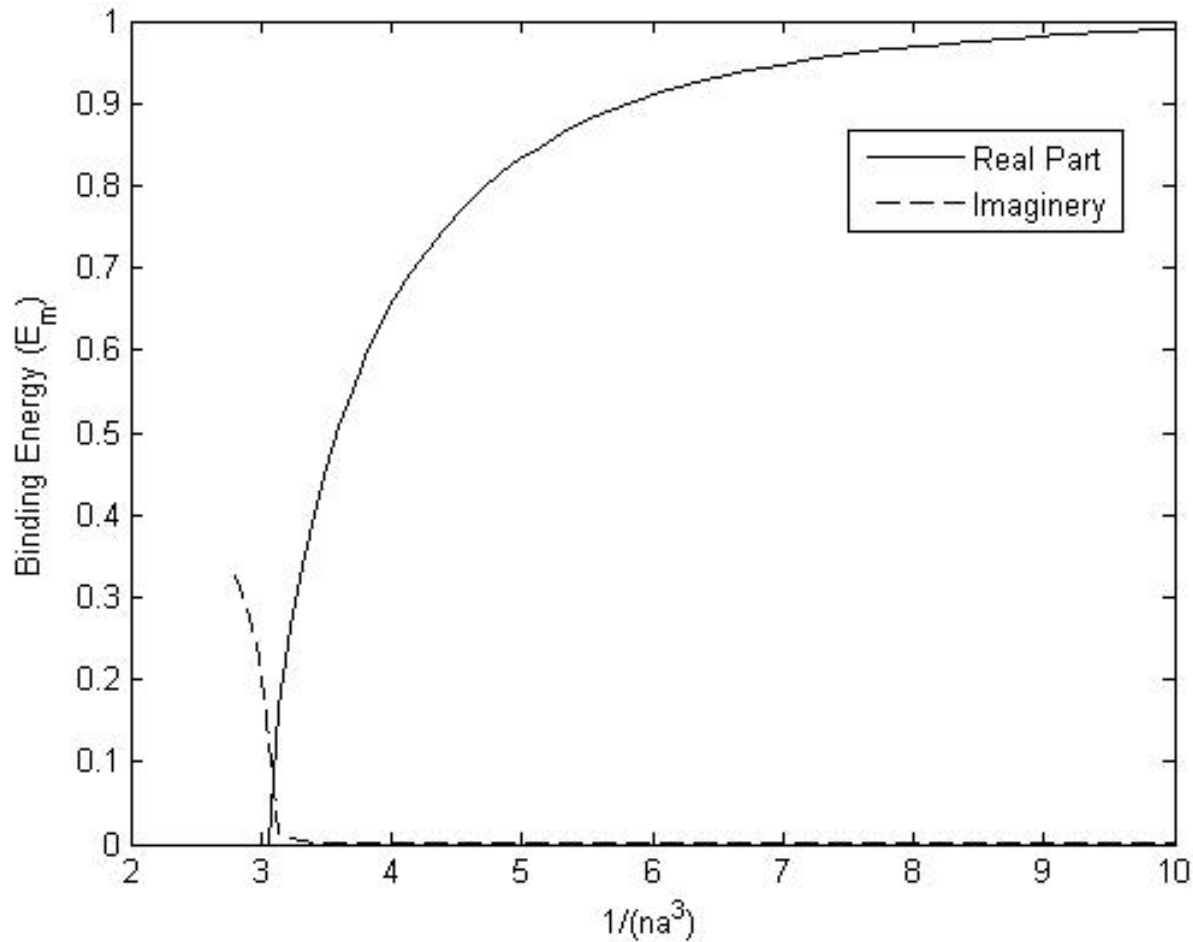


Correlation function $\chi_{ij}(\mathbf{k}, \omega) \equiv \langle T \{ d_i(\mathbf{k}, \omega) d_j(-\mathbf{k}, -\omega) \} \rangle$

$$d_1(\mathbf{k}) \equiv \sum_{\mathbf{k}'} \psi_{\mathbf{k}'} \psi_{\mathbf{k}-\mathbf{k}'}, \quad d_2(\mathbf{k}) \equiv d_1^+(-\mathbf{k}), \quad d_3(\mathbf{k}) \equiv \sum_{\mathbf{k}'} \psi_{-\mathbf{k}'}^+ \psi_{\mathbf{k}-\mathbf{k}'}$$

Dispersion Equation $\det |\chi(\mathbf{k}, \omega)| = 0$

The molecule binding energy



$$E_m \equiv \frac{\hbar^2}{ma^2}$$

When $na^3 > 0.035$, the excitation energy is purely **imaginary** and the atomic condensation state is **unstable**.

3. The molecular condensation state

Hamiltonian $H = -\psi^\dagger \frac{\hbar^2 \nabla^2}{2m} \psi + \frac{g}{2} \psi^\dagger \psi^\dagger \psi \psi, \quad g \equiv 4\pi \frac{\hbar^2 a}{m}$

Pair condensation $\Delta \equiv \frac{g}{V} \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}} \psi_{-\mathbf{k}} \rangle$

Mean-field Hamiltonian

$$H_2 = \sum_{\mathbf{k}} (\varepsilon_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{1}{2} \sum_{\mathbf{k}} (\Delta \psi_{-\mathbf{k}}^\dagger \psi_{\mathbf{k}}^\dagger + \Delta^* \psi_{\mathbf{k}} \psi_{-\mathbf{k}})$$

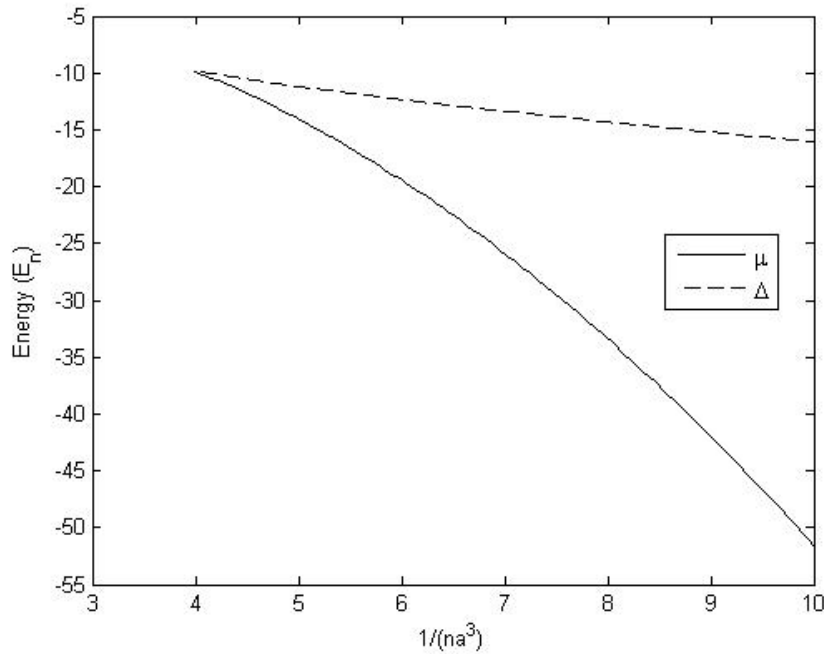
Bogoliubov Transformation $b_{\mathbf{k}} = u_{\mathbf{k}} \psi_{\mathbf{k}} - v_{\mathbf{k}} \psi_{-\mathbf{k}}^\dagger$

$$H_2 = C + \sum_{\mathbf{k}} E_{\mathbf{k}} b_{\mathbf{k}}^\dagger b_{\mathbf{k}}$$

Quasi-particle energy $E_{\mathbf{k}} = \sqrt{(\varepsilon_{\mathbf{k}} - \mu)^2 - |\Delta|^2}$

Number Equation $N = \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} \rangle$

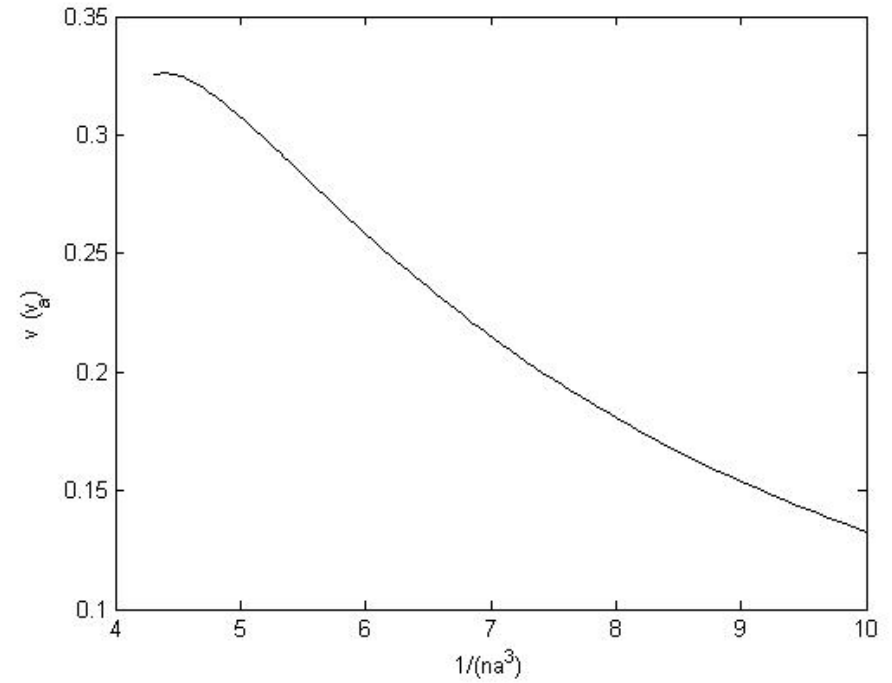
The gap



$$E_n \equiv \frac{\hbar^2 n^{2/3}}{m}$$

The solution only exists
when $1/(n^{1/3}a) > 4$.

The molecule velocity



$$v_a \equiv \frac{\hbar}{ma}$$

Molecule energy
near $k = 0$,

$$\Omega_k \approx vk$$

4. Coherent mixture of atoms and molecules

Hamiltonian $H = -\psi^\dagger \frac{\hbar^2 \nabla^2}{2m} \psi + \frac{g}{2} \psi^\dagger \psi^\dagger \psi \psi, \quad g \equiv 4\pi \frac{\hbar^2 a}{m}$

Atom and molecule condensation $\Delta \equiv \frac{g}{V} [\psi_0^2 + \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}} \psi_{-\mathbf{k}} \rangle]$

Mean-field Hamiltonian

$$H_2 = \sum_{\mathbf{k}} (\varepsilon_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{1}{2} \sum_{\mathbf{k}} (\Delta \psi_{-\mathbf{k}}^\dagger \psi_{\mathbf{k}}^\dagger + \Delta^* \psi_{\mathbf{k}} \psi_{-\mathbf{k}})$$

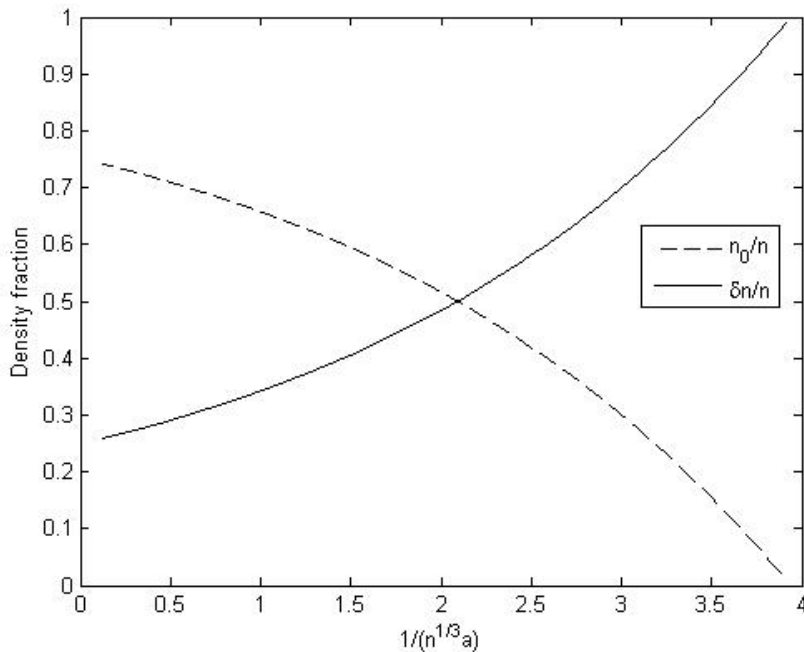
Gapless condition $|\Delta| = |\mu|$

Bogoliubov Transformation $b_{\mathbf{k}} = u_{\mathbf{k}} \psi_{\mathbf{k}} - v_{\mathbf{k}} \psi_{-\mathbf{k}}^\dagger$

$$H_2 = C + \sum_{\mathbf{k}} E_{\mathbf{k}} b_{\mathbf{k}}^\dagger b_{\mathbf{k}}, \quad E_{\mathbf{k}} = \sqrt{\varepsilon_{\mathbf{k}} (\varepsilon_{\mathbf{k}} + 2|\Delta|)}$$

Number Equation $N = N_0 + \sum_{\mathbf{k}} \langle \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} \rangle$

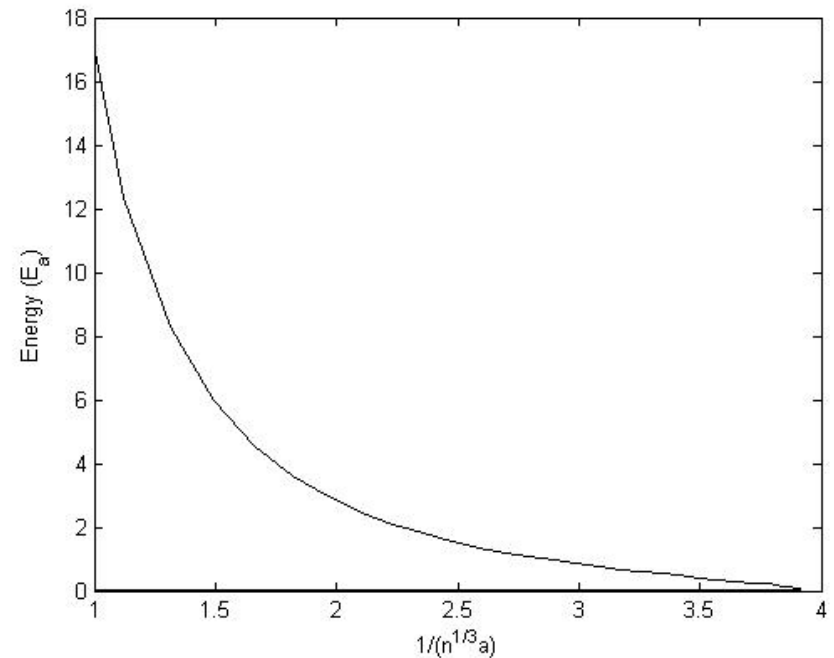
Density fraction



At resonance, the condensate density is $\frac{3}{4}$ of total density.

The signs of ψ_0^2 and $\langle \psi_{\mathbf{k}} \psi_{-\mathbf{k}} \rangle$ are opposite.

Molecule Energy

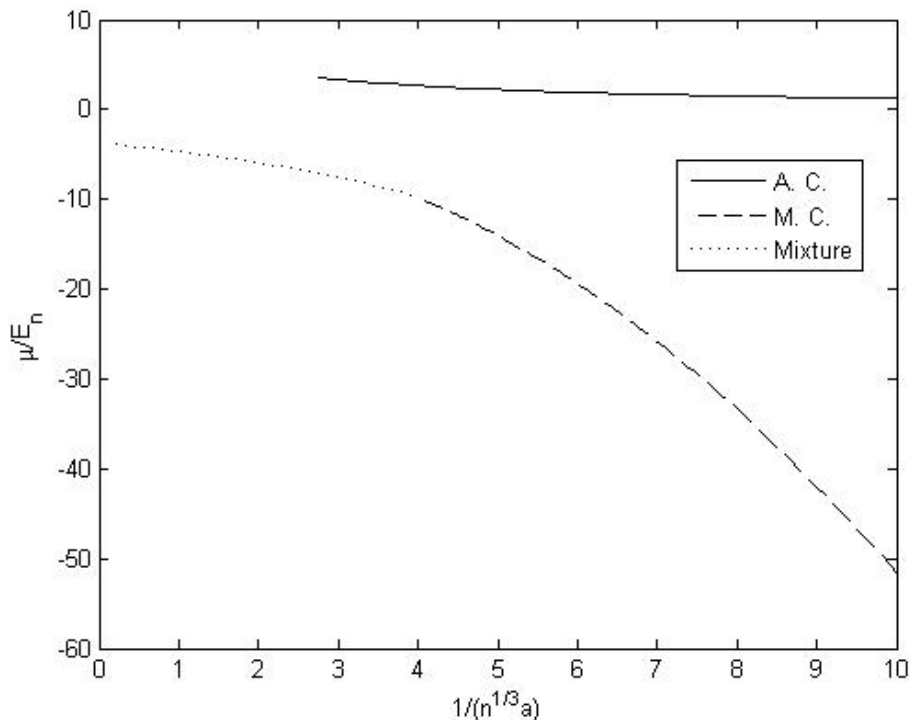


$$E_a \equiv \frac{\hbar^2}{ma^2}$$

Molecular excitation opens a gap.

5. Conclusion

We have studied three possible phases of a repulsive Bose gas. The atomic condensation state is metastable when $na^3 < 0.035$. There is a transition between the molecular condensation and the mixture states when $na^3 = 0.015$. The atomic excitation is gapless in the AC and mixture states. The molecular excitation is gapless in the MC state.



$$E_n \equiv \frac{\hbar^2 n^{2/3}}{m}$$